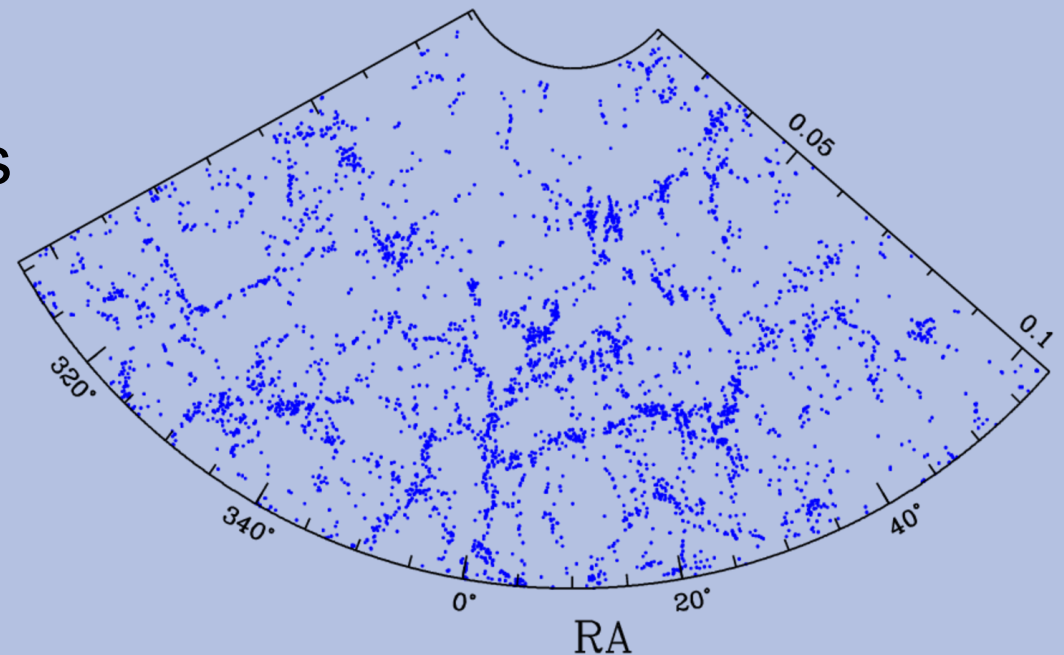
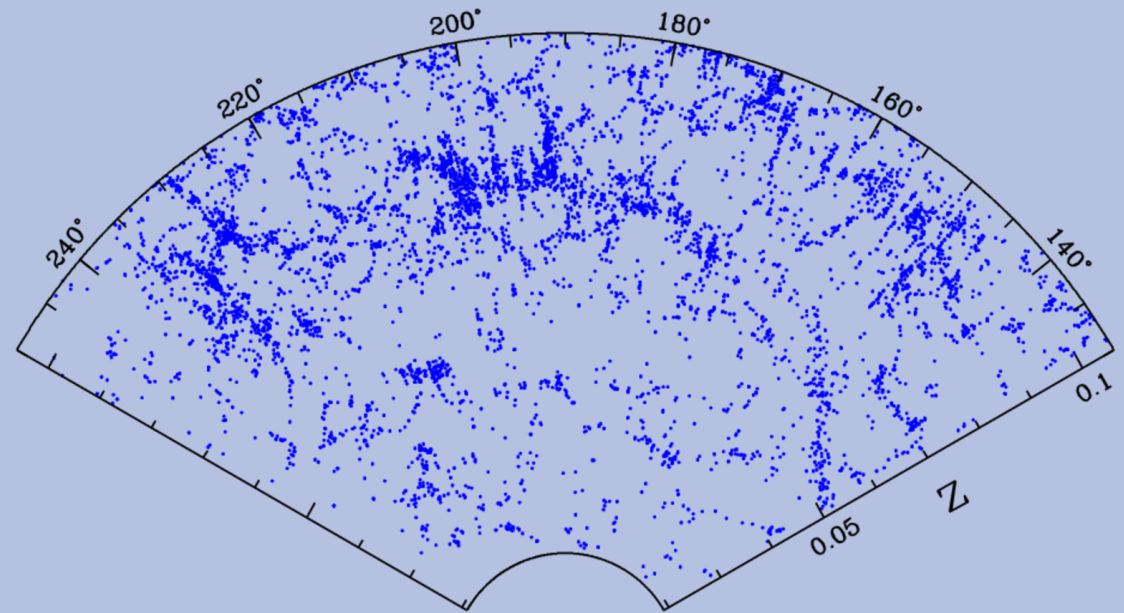


# Statistics of the Galaxy Distribution

measure the environment  
around individual galaxies

or

measure average statistics  
for a sample of galaxies



# The 2-point Correlation Function

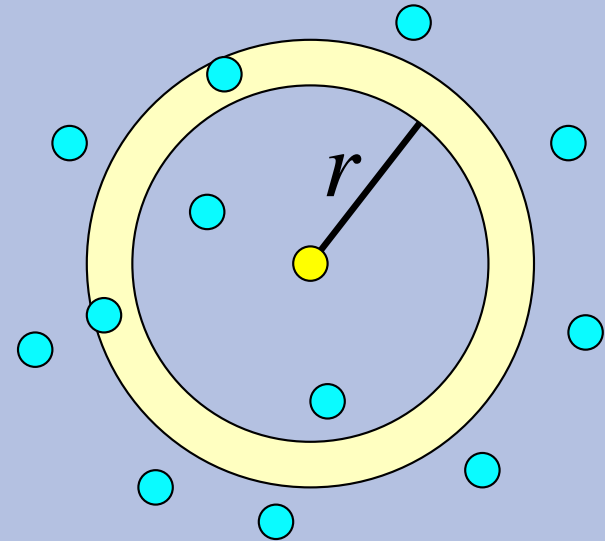
The excess probability that two galaxies are separated by a distance  $r$  relative to that for a random distribution.

For a random point distribution of number density  $n$ , the number of points at a distance between  $r$  and  $r+dr$  from any one point is:

$$n \cdot dV = n \cdot 4\pi r^2 dr$$

The total number density of pairs at this separation is then:

$$\frac{n}{2} \times n \cdot 4\pi r^2 dr = n^2 2\pi r^2 dr$$



# The 2-point Correlation Function

For a point distribution that is not random, the number density of pairs  
At this separation is:

$$n^2 2\pi r^2 dr [1 + \xi(r)]$$

The correlation function is then equal to:

$$\xi(r) = \frac{n_{\text{pairs, data}}(r)}{n_{\text{pairs, rand}}(r)} - 1$$

$$\xi(r) = \frac{DD(r)}{RR(r)} - 1$$

# The 2-point Correlation Function

## Complications

- When the sample volume is complex, calculate number of random pairs using an actual generated random data set that occupies the same volume as the data.
- In practice, we often use other estimators. For example, the most commonly used estimator for galaxy samples is the Landy-Szalay estimator:

$$\xi(r) = \frac{DD - 2DR + RR}{RR}$$

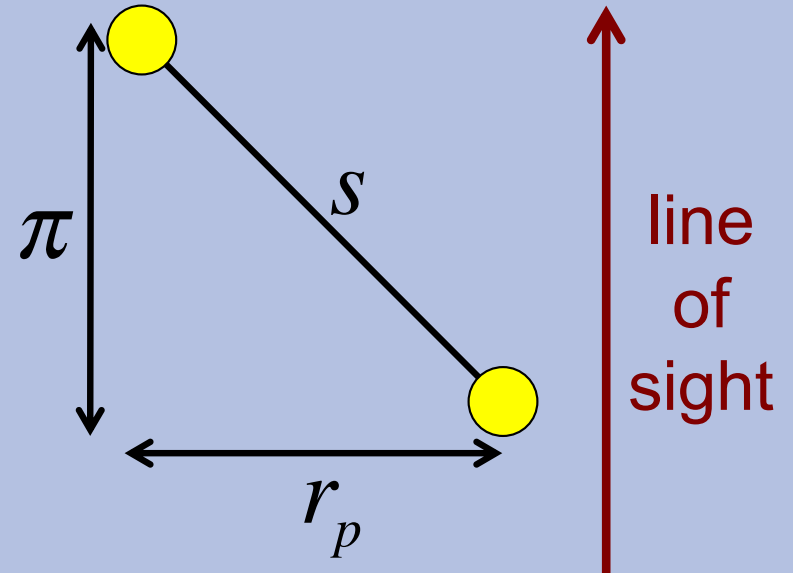
- Must normalize the DD, DR, and RR terms when the number of data and random points are not the same.

# The 2-point Correlation Function

## Complications

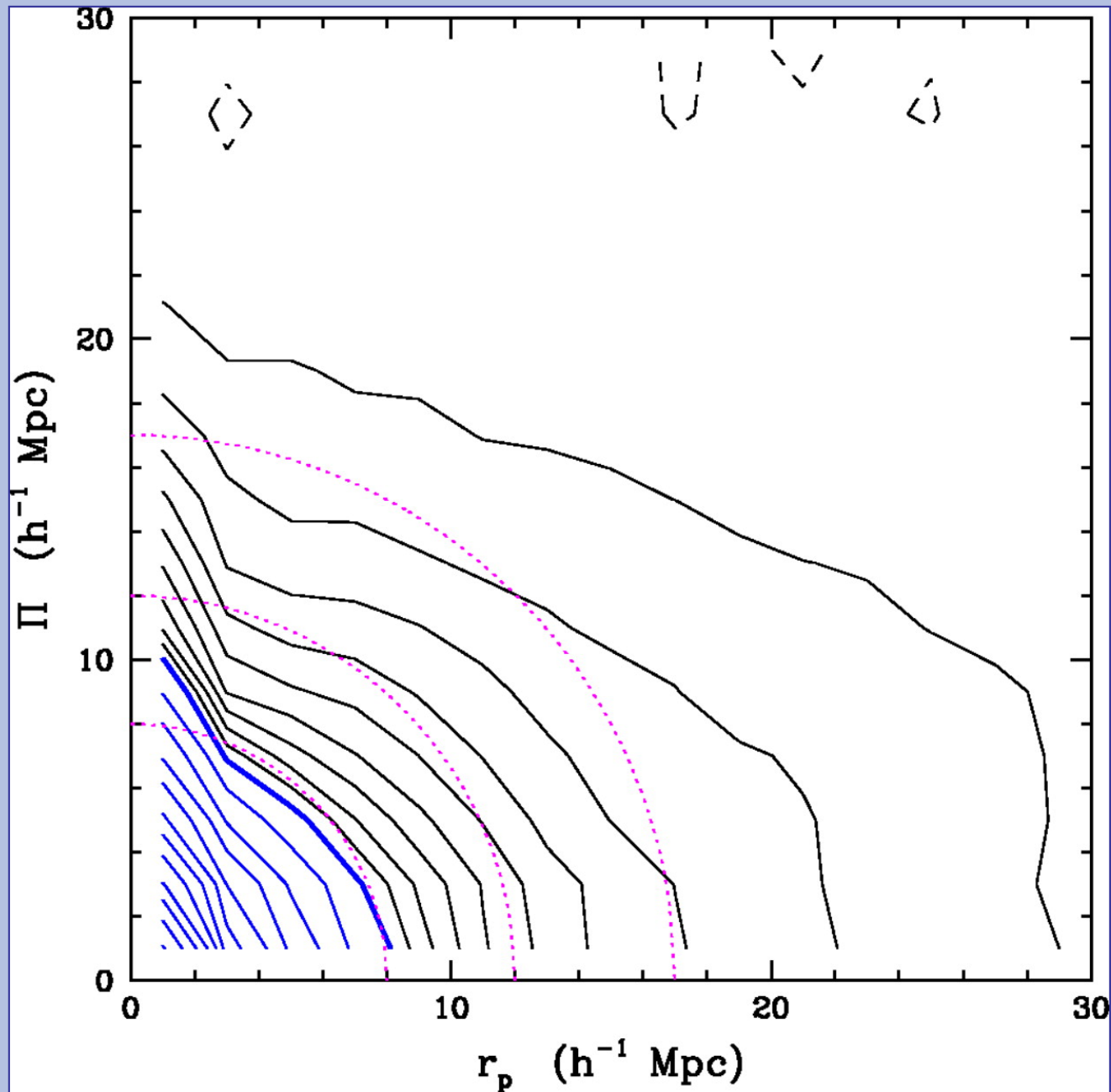
- To deal with redshift distortions, we usually measure a *projected* correlation function.

$$\xi(r_p, \pi)$$



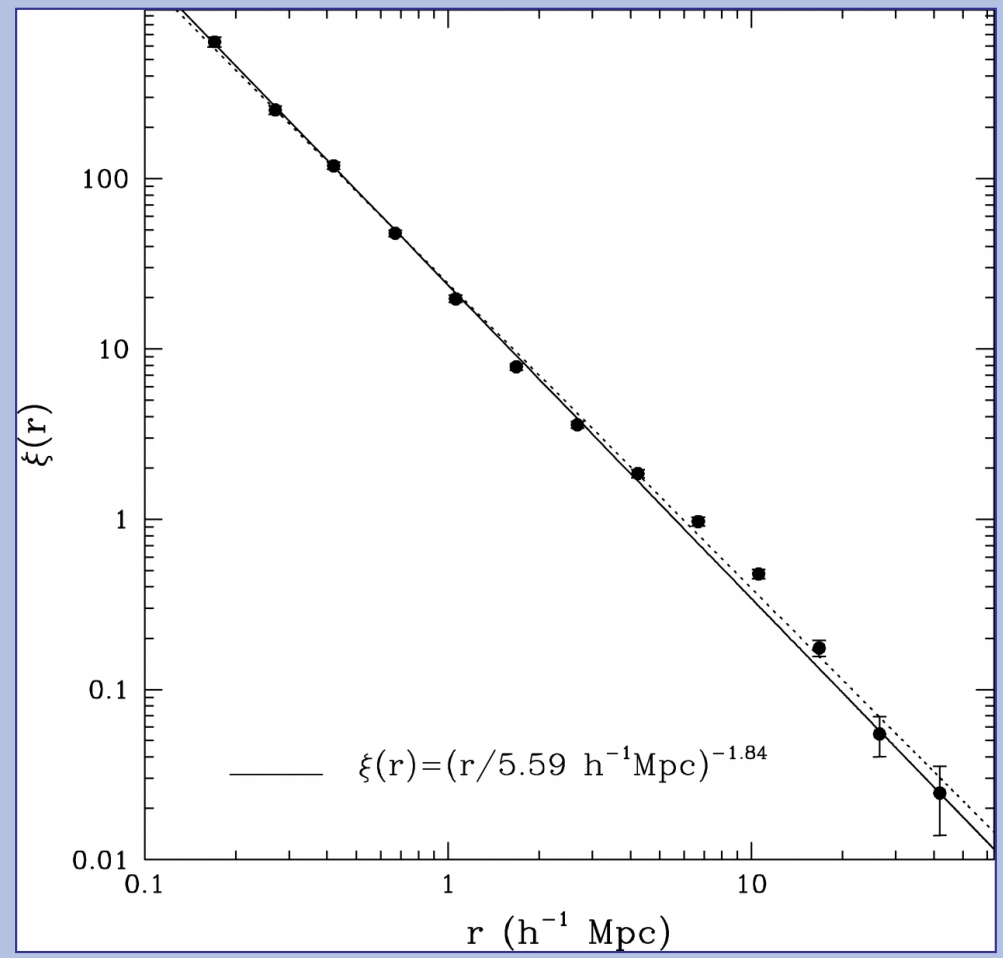
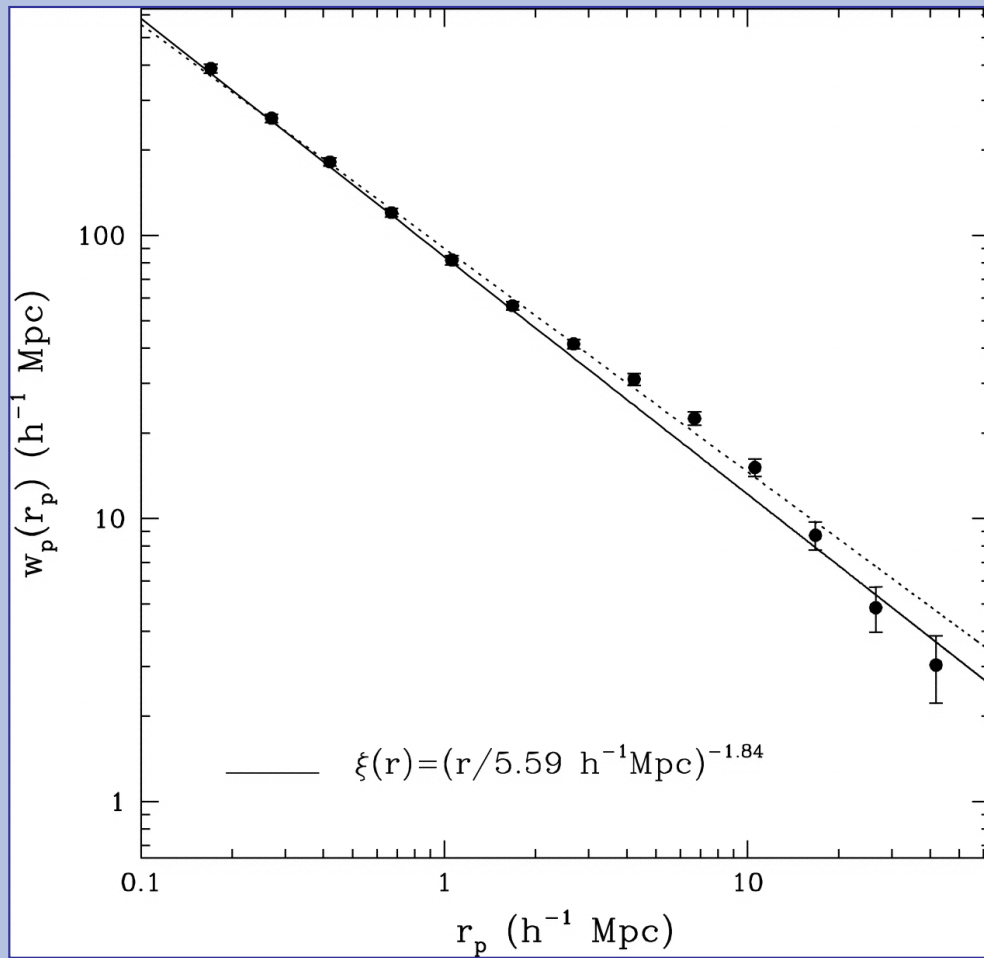
$$w_p(r_p) = 2 \int_0^{\pi_{\max}} \xi(r_p, \pi) d\pi$$

# The 2-point Correlation Function



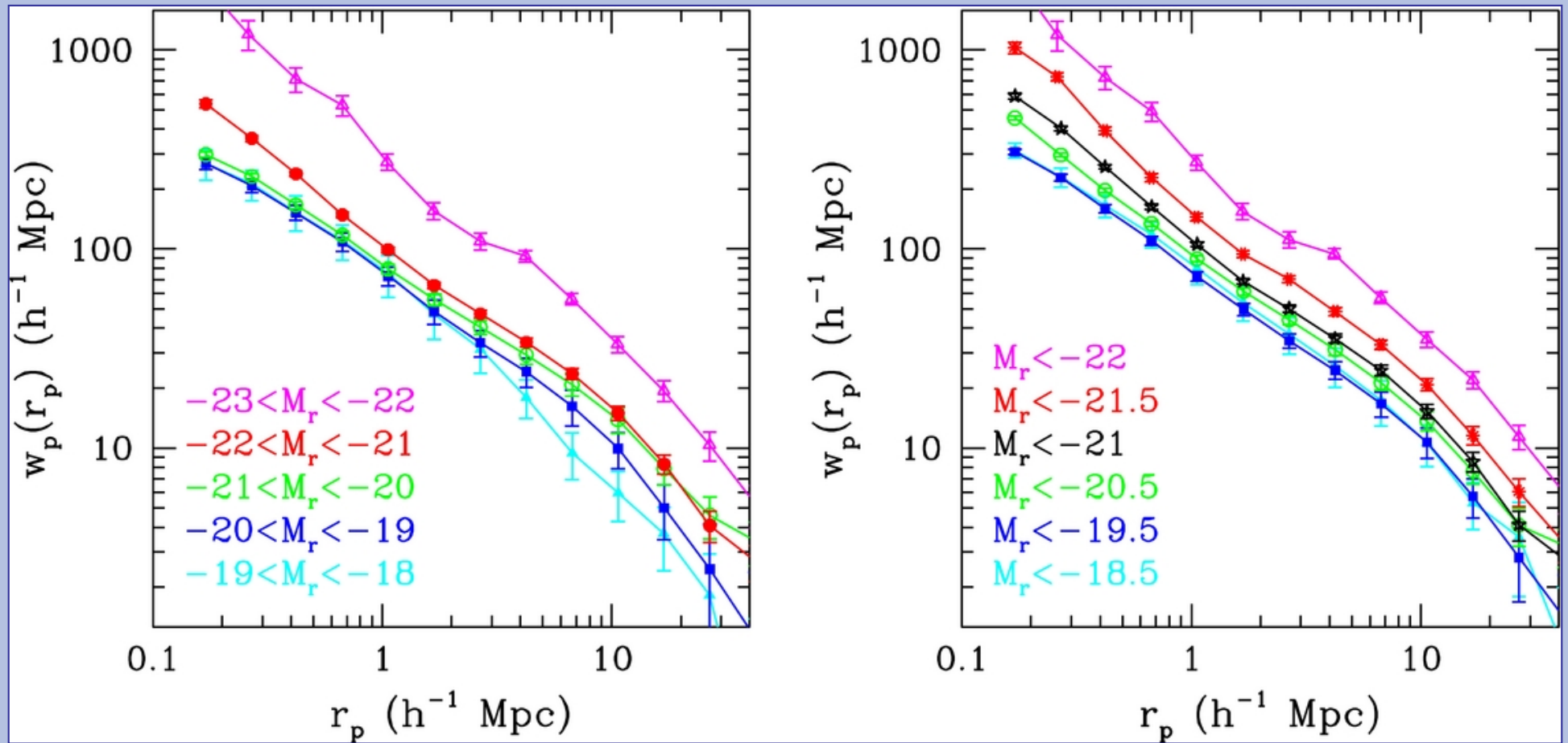
Zehavi et al. (2005)

# The 2-point Correlation Function



Zehavi et al. (2005)

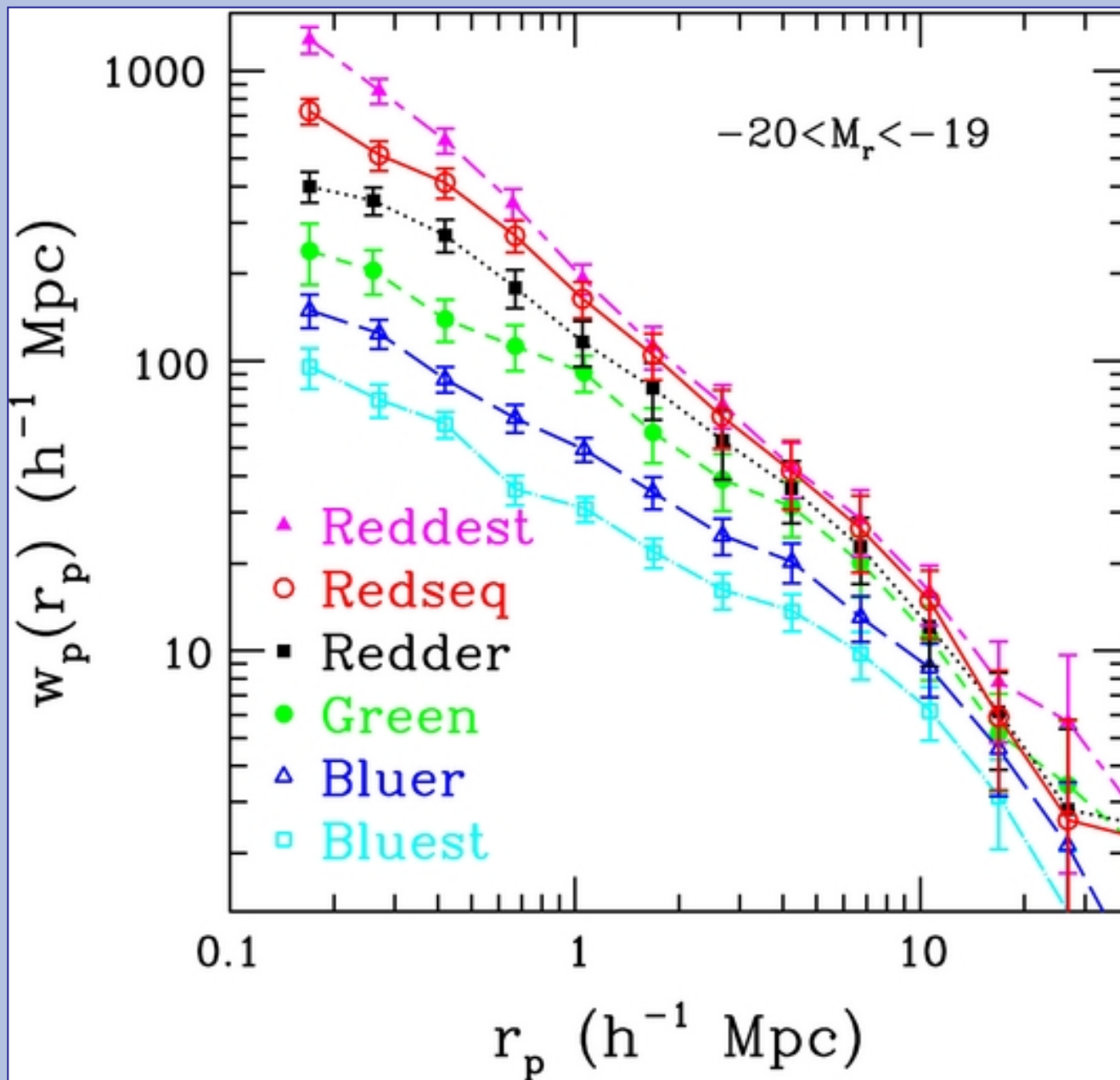
# The 2-point Correlation Function



Zehavi et al. (2011)

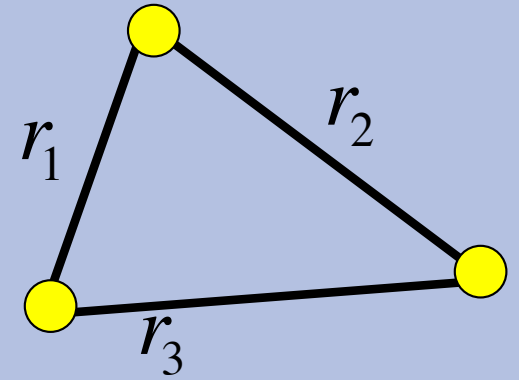
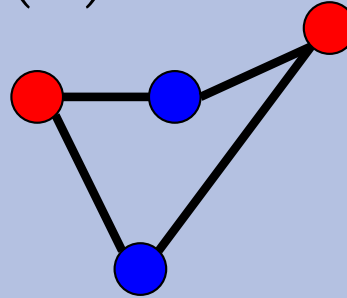


# The 2-point Correlation Function

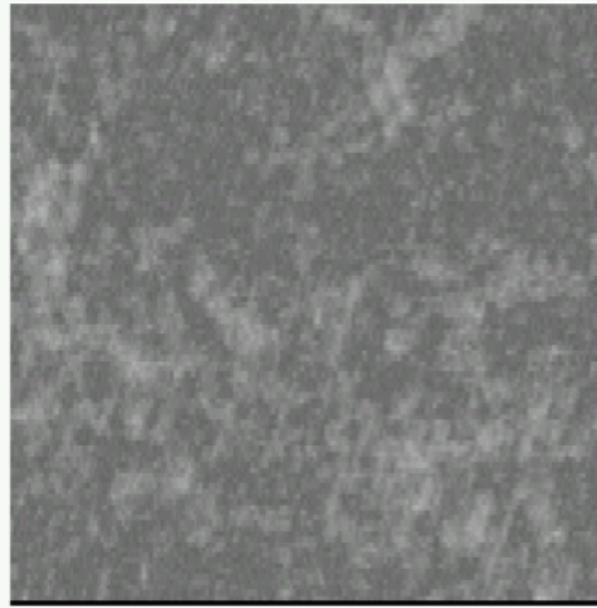
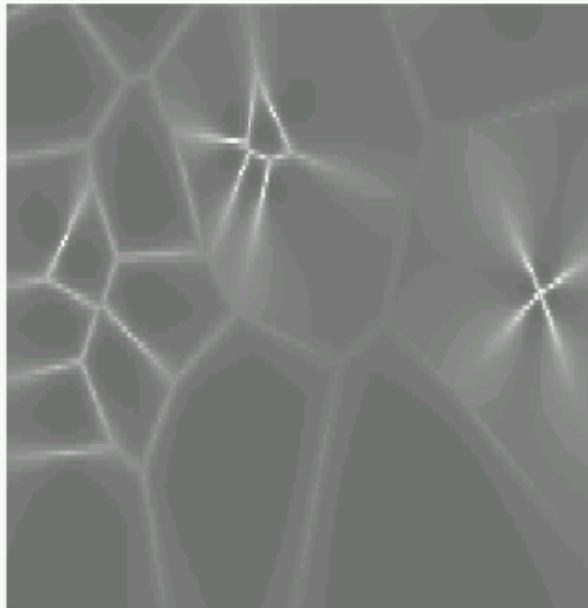


# Other Correlation Functions

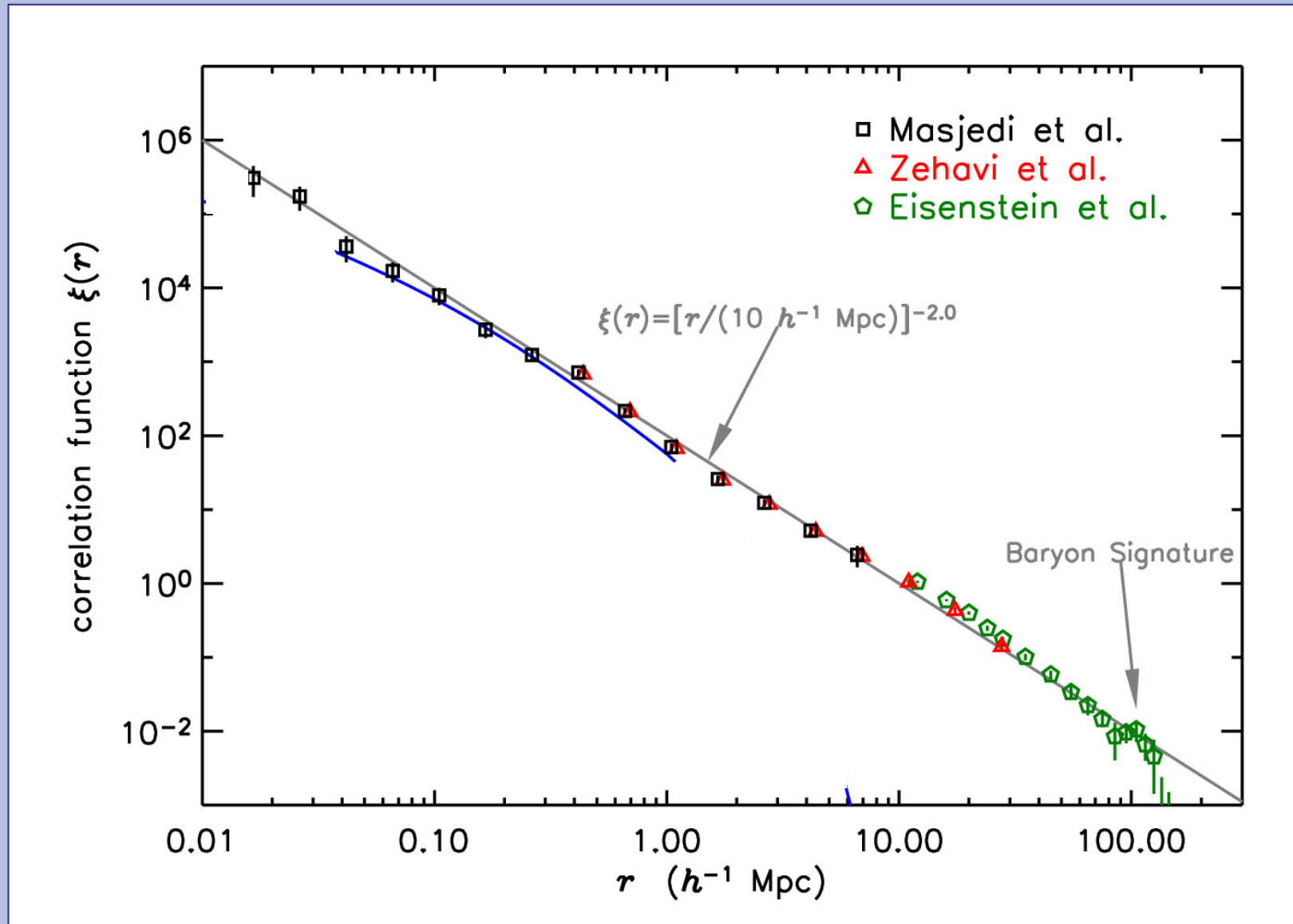
- Angular correlation function  $\omega(\theta)$
- Cross-correlation function
- Higher order: three-point, etc



• same 2PCF but very different distributions



# The 2-point Correlation Function



Masjedi et al. (2006)

# The 2-point Correlation Function

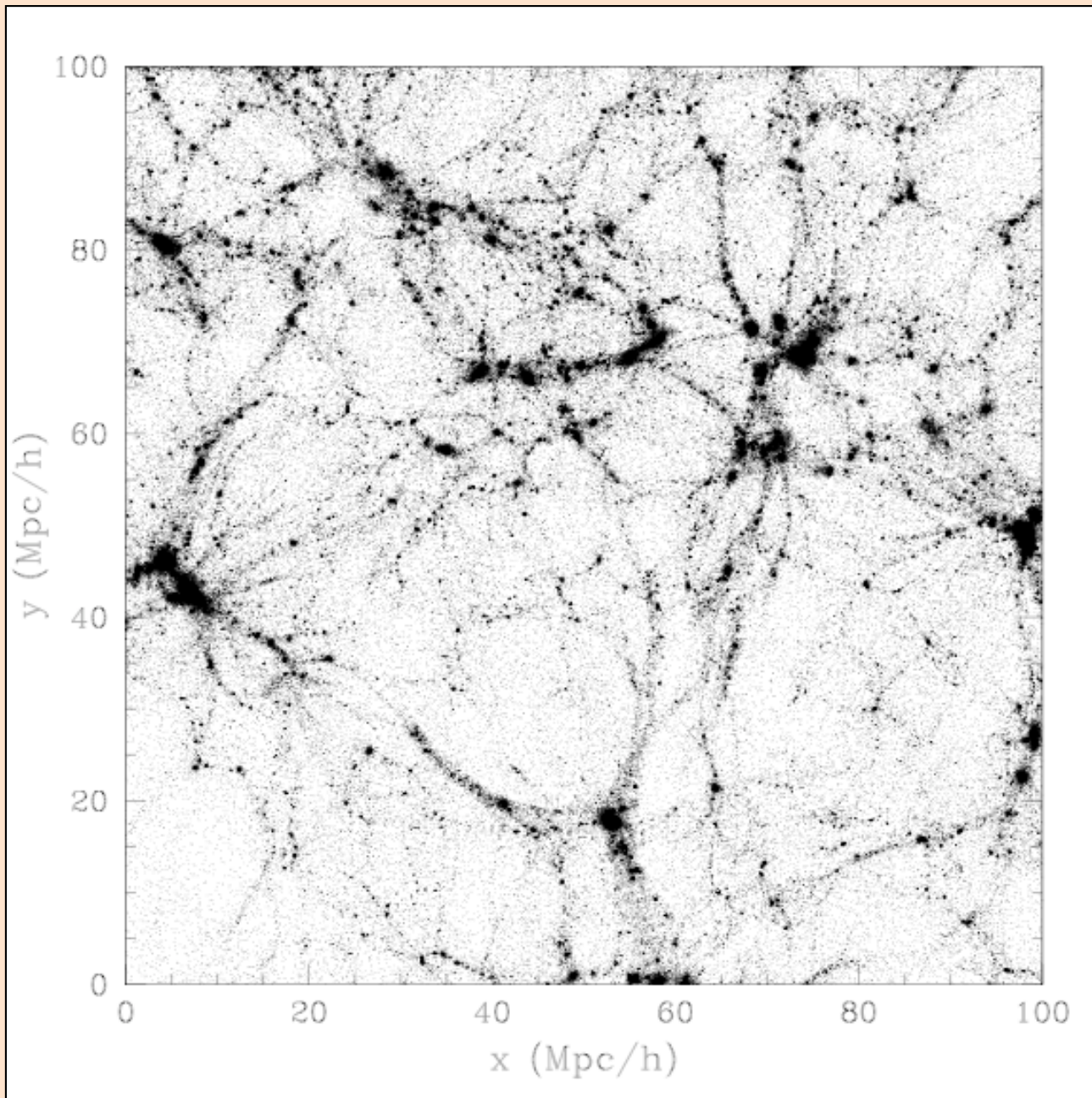
Galaxy “bias” refers to the amount of galaxy clustering relative to the clustering of the underlying dark matter

$$b = \sqrt{\frac{\xi_{\text{galaxy}}}{\xi_{\text{mass}}}}$$

This is a function of scale  $r$ , but on large scales it becomes constant.

For example, the bias of Milky Way – like galaxies is  $b \sim 1$   
the bias of Luminous Red Galaxies is  $b \sim 2$

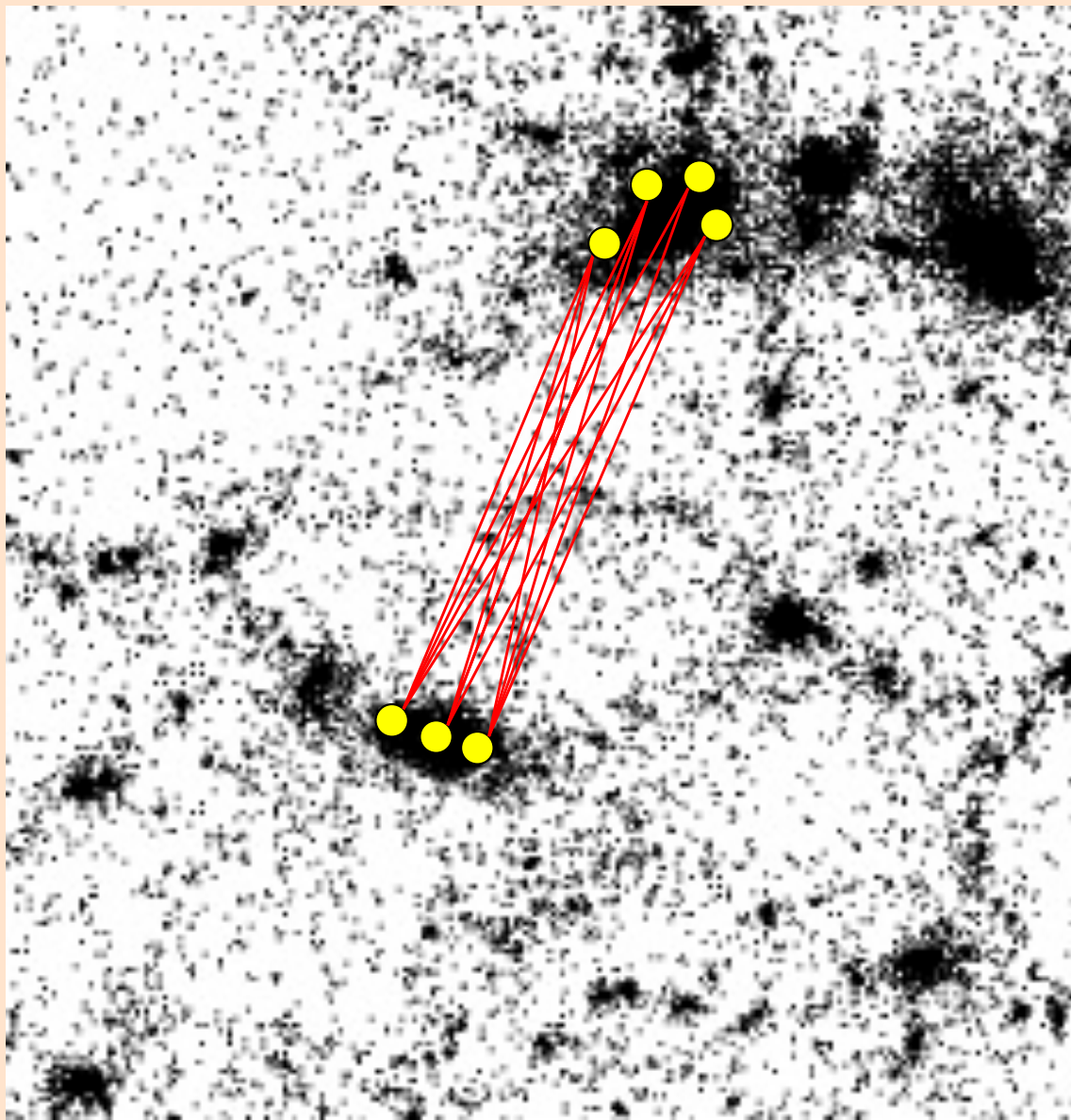
# Power Spectrum



Density field

$$\delta(\vec{x}) = \frac{\rho(\vec{x}) - \bar{\rho}}{\bar{\rho}}$$

# Power Spectrum



$$\xi(r) = \frac{DD}{RR} - 1$$

$$DD = \rho(\vec{x})\rho(\vec{x} + \mathbf{r})$$

$$RR = \bar{\rho}\bar{\rho}$$

# Power Spectrum

$$\xi(r) = \frac{DD}{RR} - 1 \quad DD = \rho(\vec{x})\rho(\vec{x} + r) \quad RR = \bar{\rho}\bar{\rho}$$

$$\xi(r) = \left\langle \frac{\rho(\vec{x})}{\bar{\rho}} \frac{\rho(\vec{x} + r)}{\bar{\rho}} \right\rangle - 1 \quad \delta(\vec{x}) = \frac{\rho(\vec{x})}{\bar{\rho}} - 1$$

$$= \langle (\delta(\vec{x}) + 1)(\delta(\vec{x} + r) + 1) \rangle - 1$$

$$= \langle \delta(\vec{x}) \rangle + \langle \delta(\vec{x} + r) \rangle + \langle \delta(\vec{x})\delta(\vec{x} + r) \rangle$$

$$= \langle \delta(\vec{x})\delta(\vec{x} + r) \rangle$$

# Power Spectrum

## Density field

$$\delta(\vec{x}) = \frac{\rho(\vec{x}) - \bar{\rho}}{\bar{\rho}}$$

## Correlation function

$$\xi(r) = \langle \delta(\vec{x}) \delta(\vec{x} + r) \rangle$$

## Fourier density modes

$$\delta_{\vec{k}} = \int \delta(\vec{x}) e^{i\vec{k} \cdot \vec{x}} d^3 \vec{x}$$

## Power spectrum

$$P(k) = \langle |\delta_{\vec{k}}|^2 \rangle$$

$$P(k) = \int \xi(r) e^{i\vec{k} \cdot \vec{r}} d^3 \vec{r}$$



# Power Spectrum

Any density field can be decomposed into an infinite set of modes (i.e., sine waves)  $\delta_{\vec{k}}$

Each mode has a

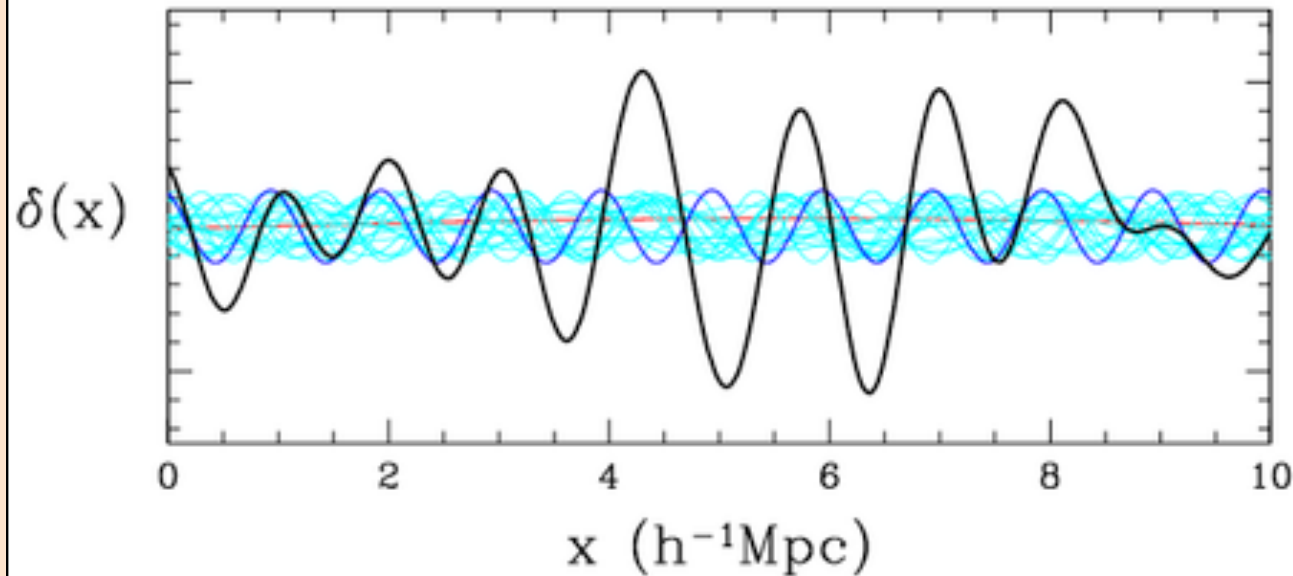
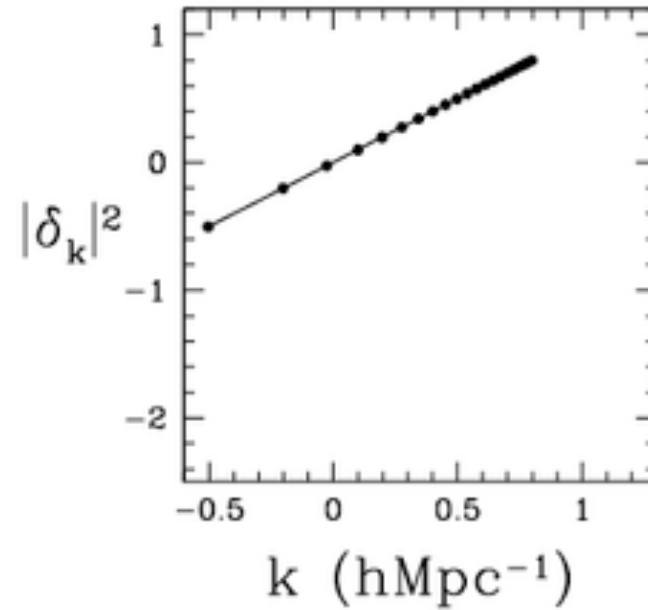
- wavelength  $\lambda$  or wavenumber  $k = \frac{2\pi}{\lambda}$
- amplitude  $|\delta_{\vec{k}}|$
- phase  $e^{-i\theta}$

The power spectrum is the amplitude as a function of k

# Power Spectrum

$$P(k) = k^1$$

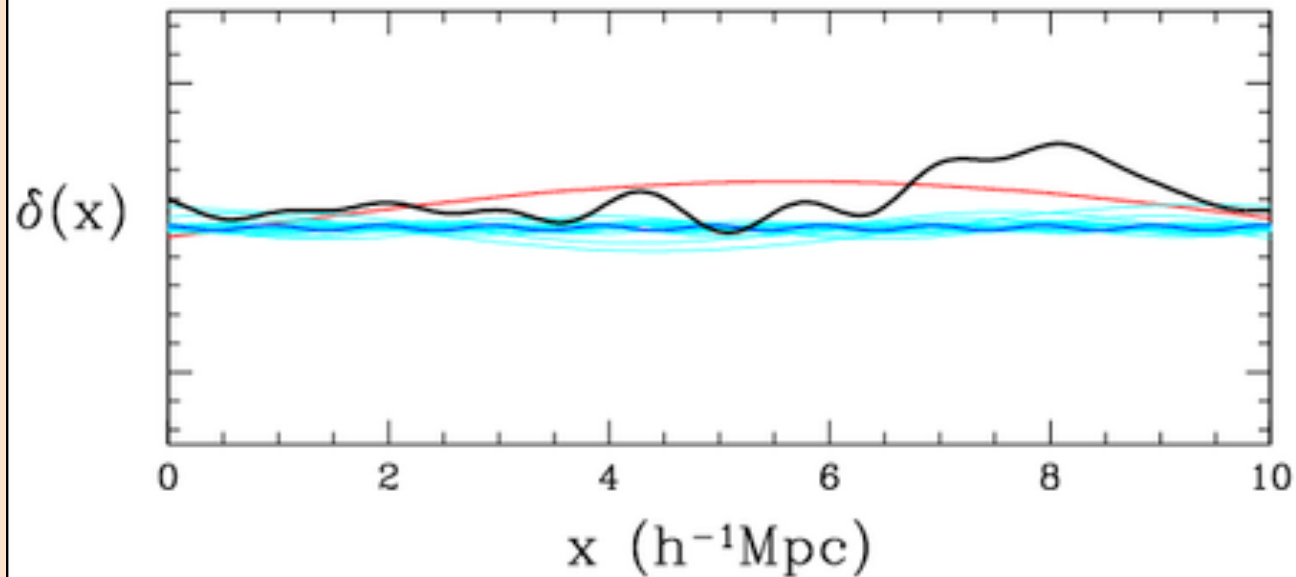
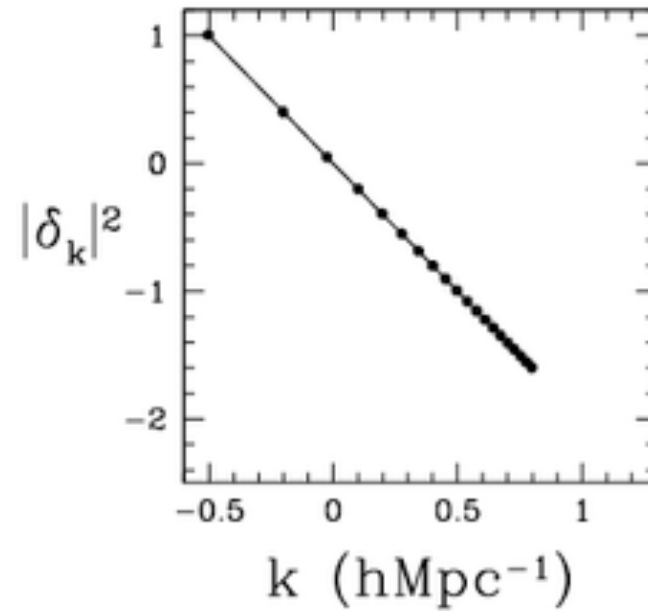
# of modes=20  
random seed=1001



# Power Spectrum

$$P(k) = k^{-2.0}$$

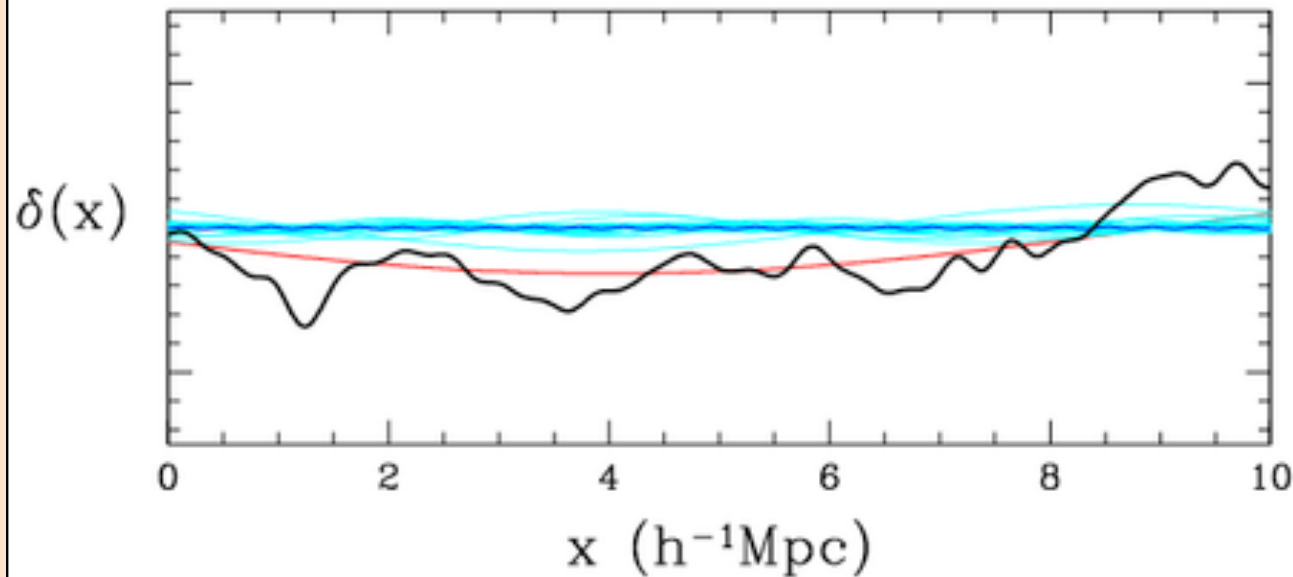
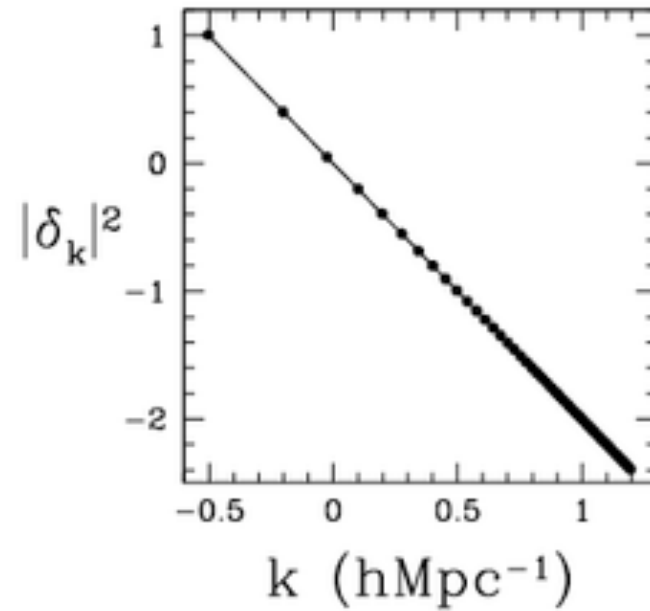
# of modes=20  
random seed=1001



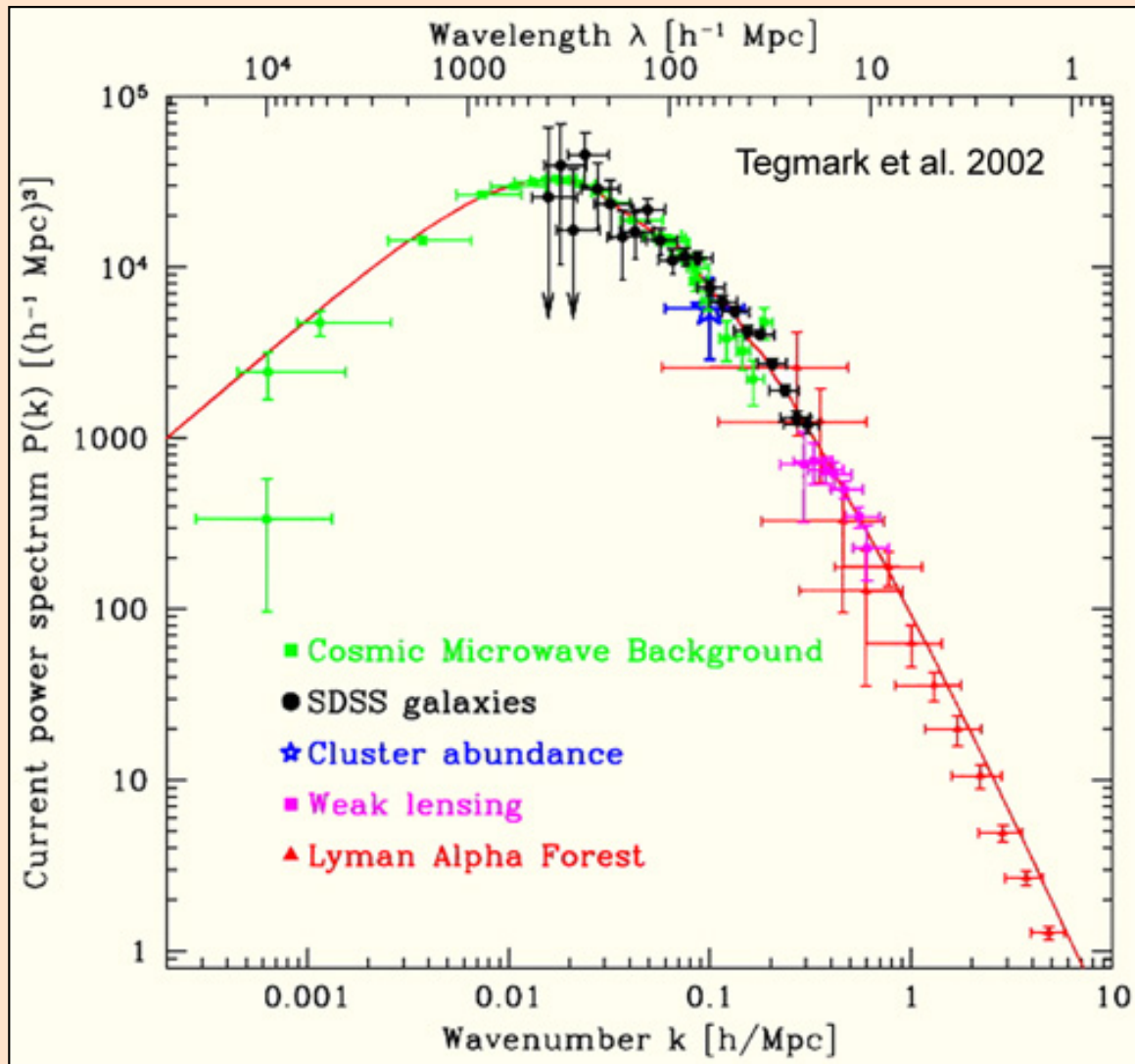
# Power Spectrum

$$P(k) = k^{-2.0}$$

# of modes=50  
random seed=19



# Power Spectrum



# Power Spectrum

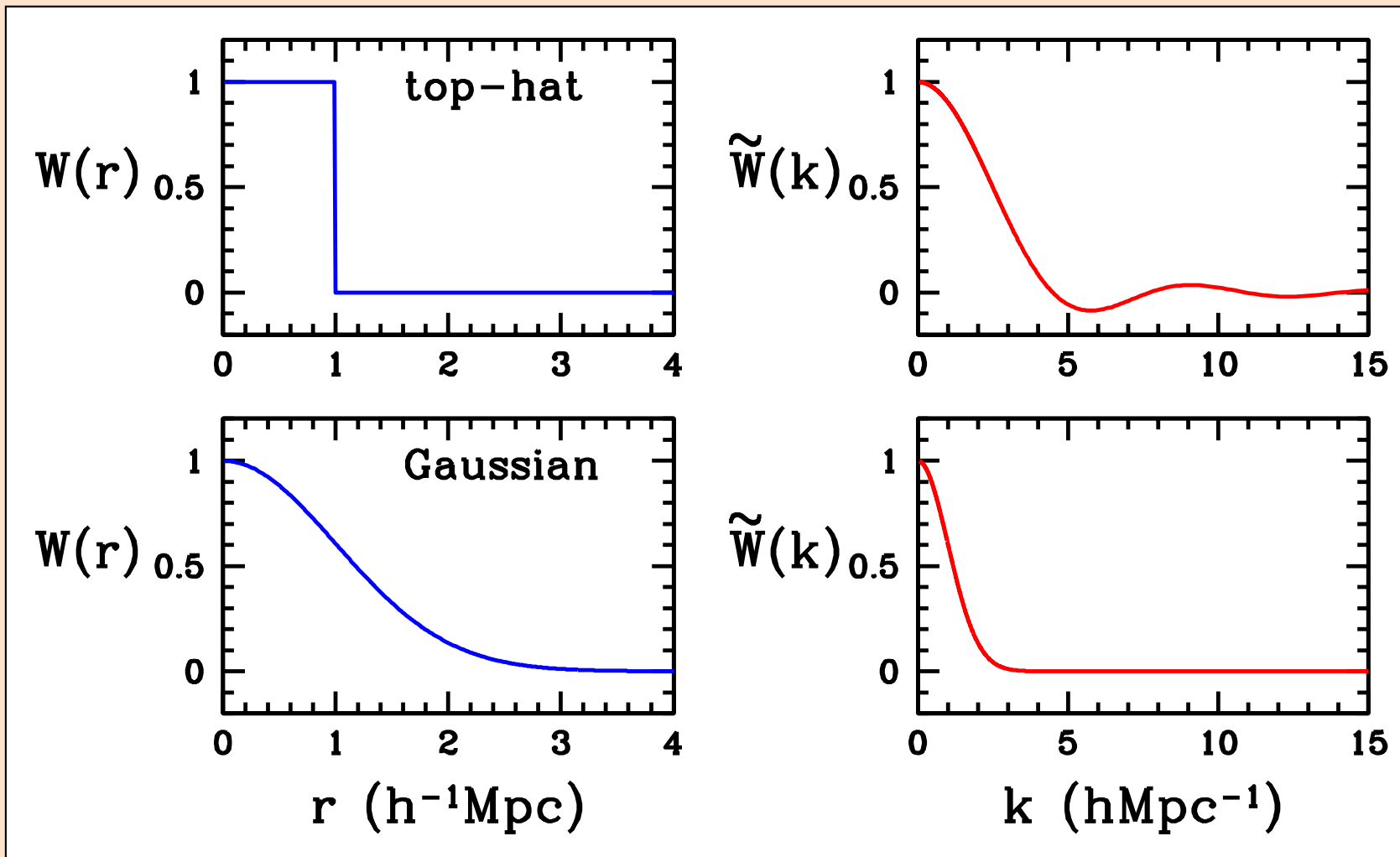
Window function: filter used to smooth density field

- Gaussian filter of scale  $R$       $W_R(r) = e^{-r^2/2R^2}$

- Top-hat filter of scale  $R$       $W_R(r) = \begin{cases} 1 & r < R \\ 0 & r > R \end{cases}$

In Fourier space:      $\tilde{W}_R(k) = \int W_R(r) e^{i\vec{k}\cdot\vec{r}} d^3r$

# Power Spectrum



# Power Spectrum

Density field, smoothed with window function

$$\delta_R(\vec{x}) = \int \delta(\vec{x}') W_R(|\vec{x}' - \vec{x}|) d^3 x'$$

Mean density, smoothed with window function

$$\bar{\delta}_R = \langle \delta_R(\vec{x}) \rangle = 0 \quad \text{since} \quad \delta = \frac{\rho - \bar{\rho}}{\bar{\rho}}$$

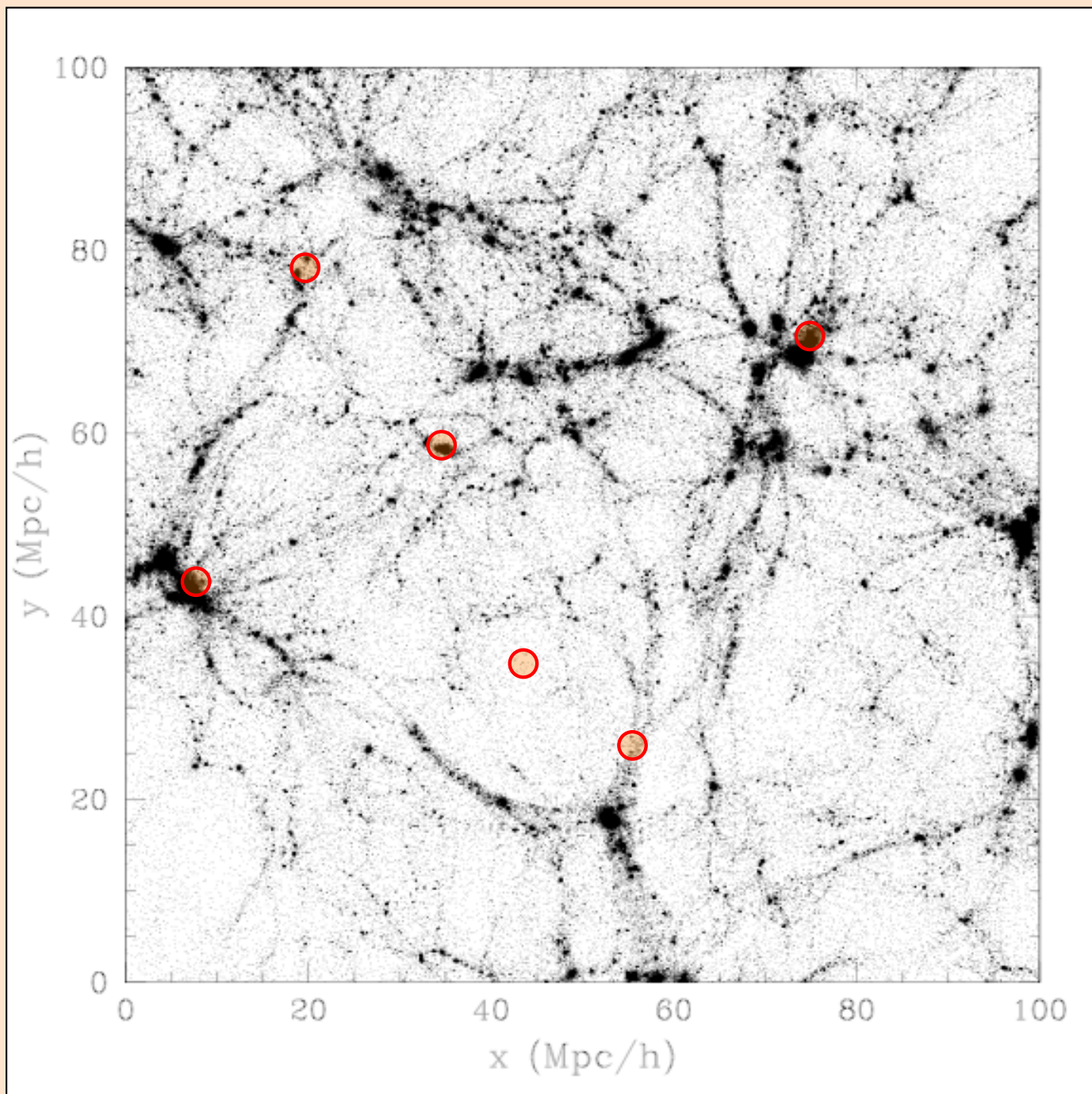
Variance of smoothed density field

$$\sigma_R^2 = \langle \delta_R(\vec{x})^2 \rangle$$

$$\sigma_R^2 = \int P(k) \tilde{W}_R^2(k) d^3 k$$



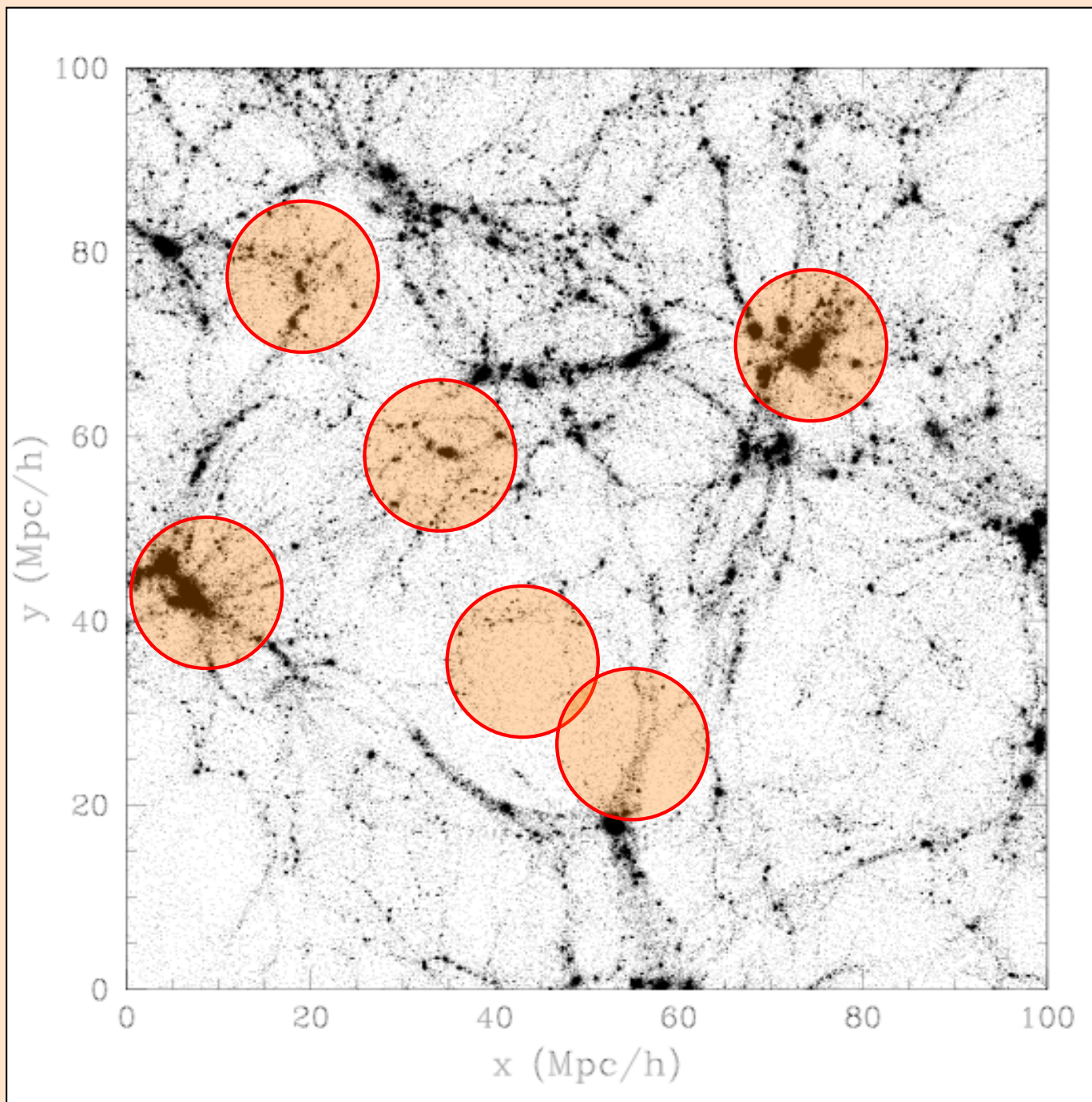
# Power Spectrum



$R=2$  Mpc/h

Top-hat filter

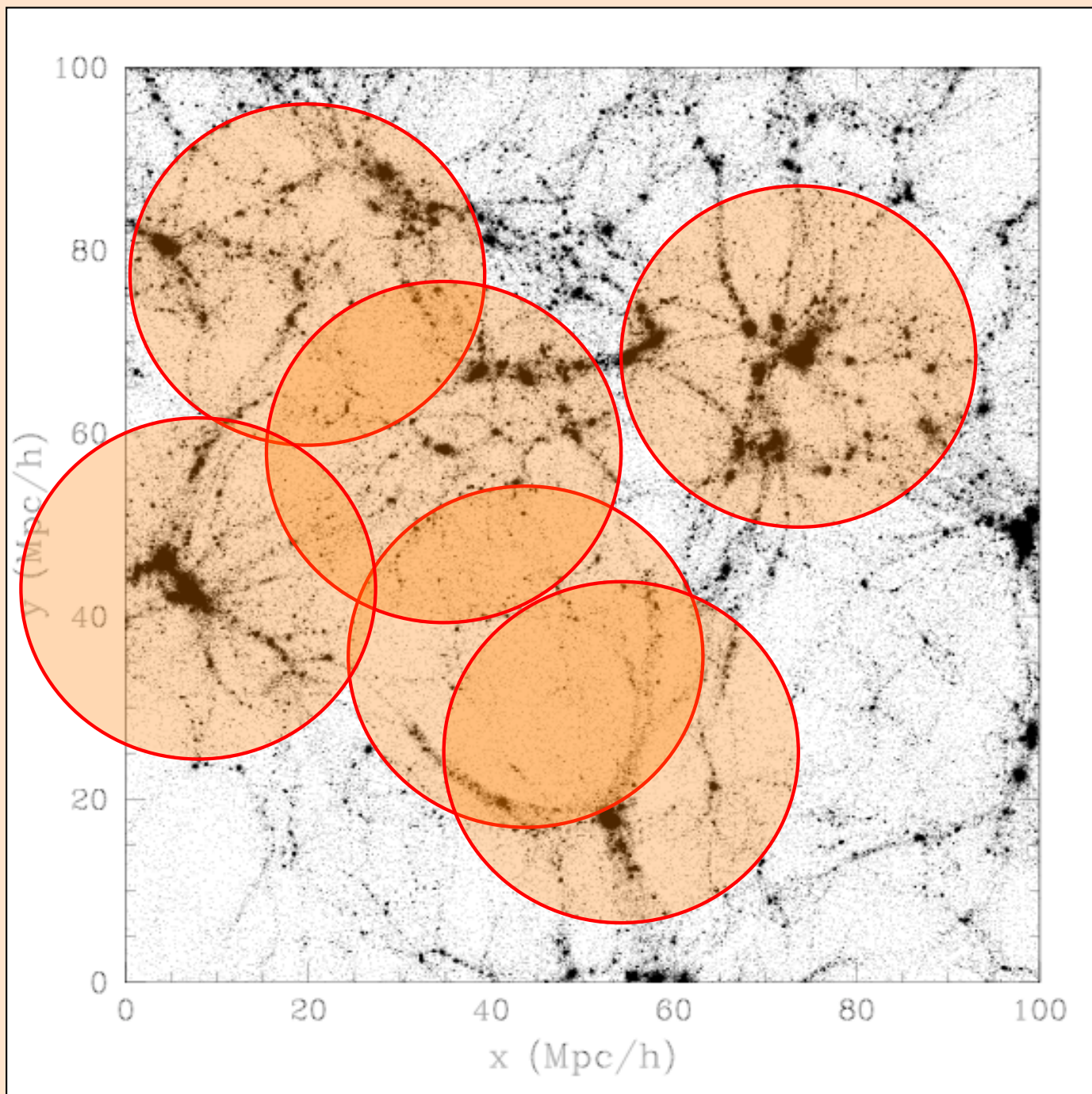
# Power Spectrum



$R=8 \text{ Mpc}/h$

Top-hat filter

# Power Spectrum



$R=20$  Mpc/h

Top-hat filter

# Power Spectrum

$$\sigma_R^2 = \frac{1}{N} \sum (\delta_R - \overline{\delta_R})^2 = \langle \delta_R^2 \rangle$$

0

The variance is large on small scales and approaches zero on large scales.

# Power Spectrum

$\sigma_R^2$  Is the variance of the matter density field

It also sets the amplitude of the matter power spectrum on scale  $R$

$$\sigma_R^2 = \langle \delta_R(\vec{x})^2 \rangle$$
$$\sigma_R^2 = \int P(k) \tilde{W}_R^2(k) d^3k$$

For a power spectrum with a power-law shape  $P(k) \sim k^n$ , defining the variance on one scale sets the amplitude on all scales. Also, any window function will do.

We choose a top-hat filter of  $R=8$  Mpc/h to describe the amplitude of  $P(k)$