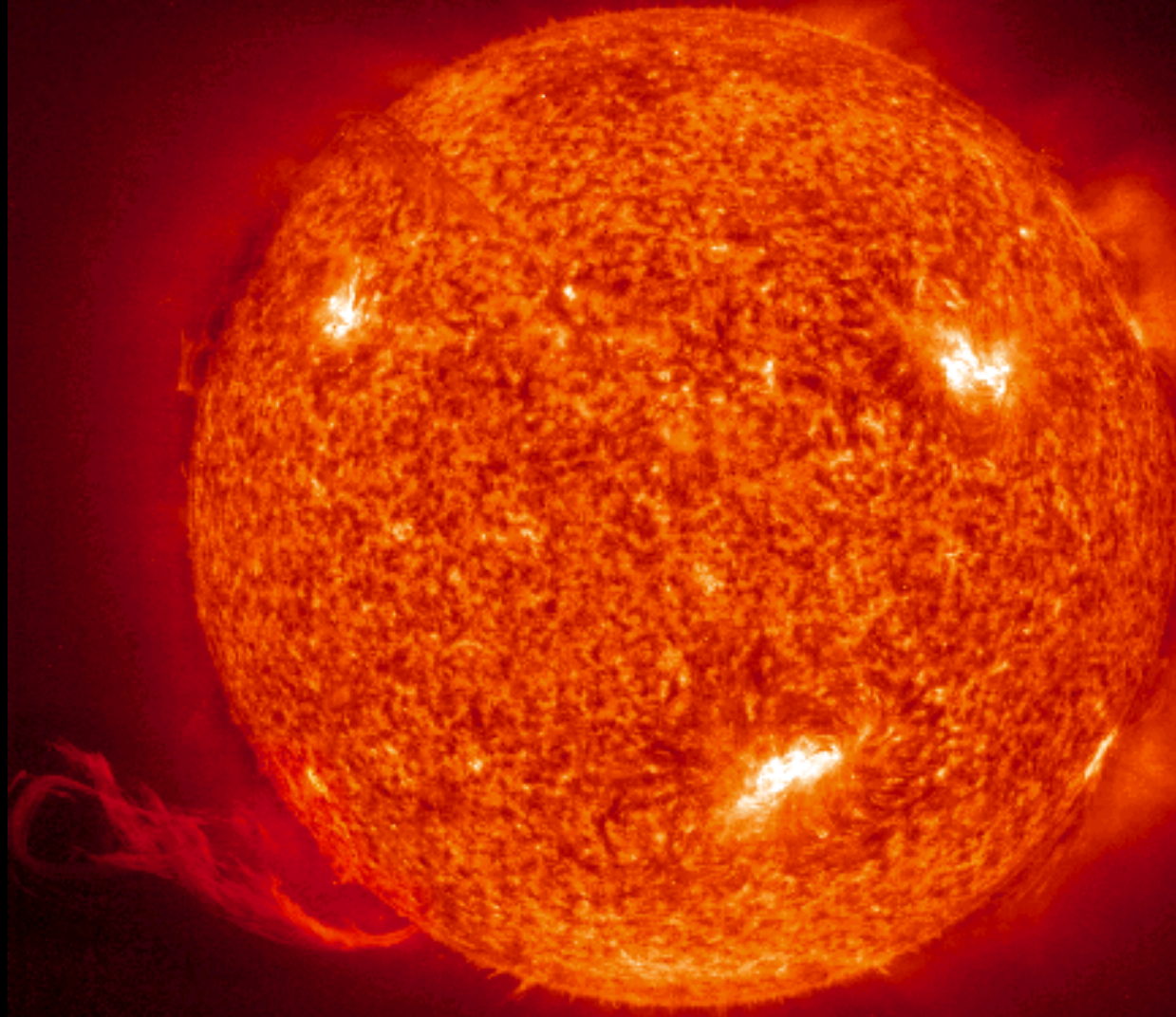


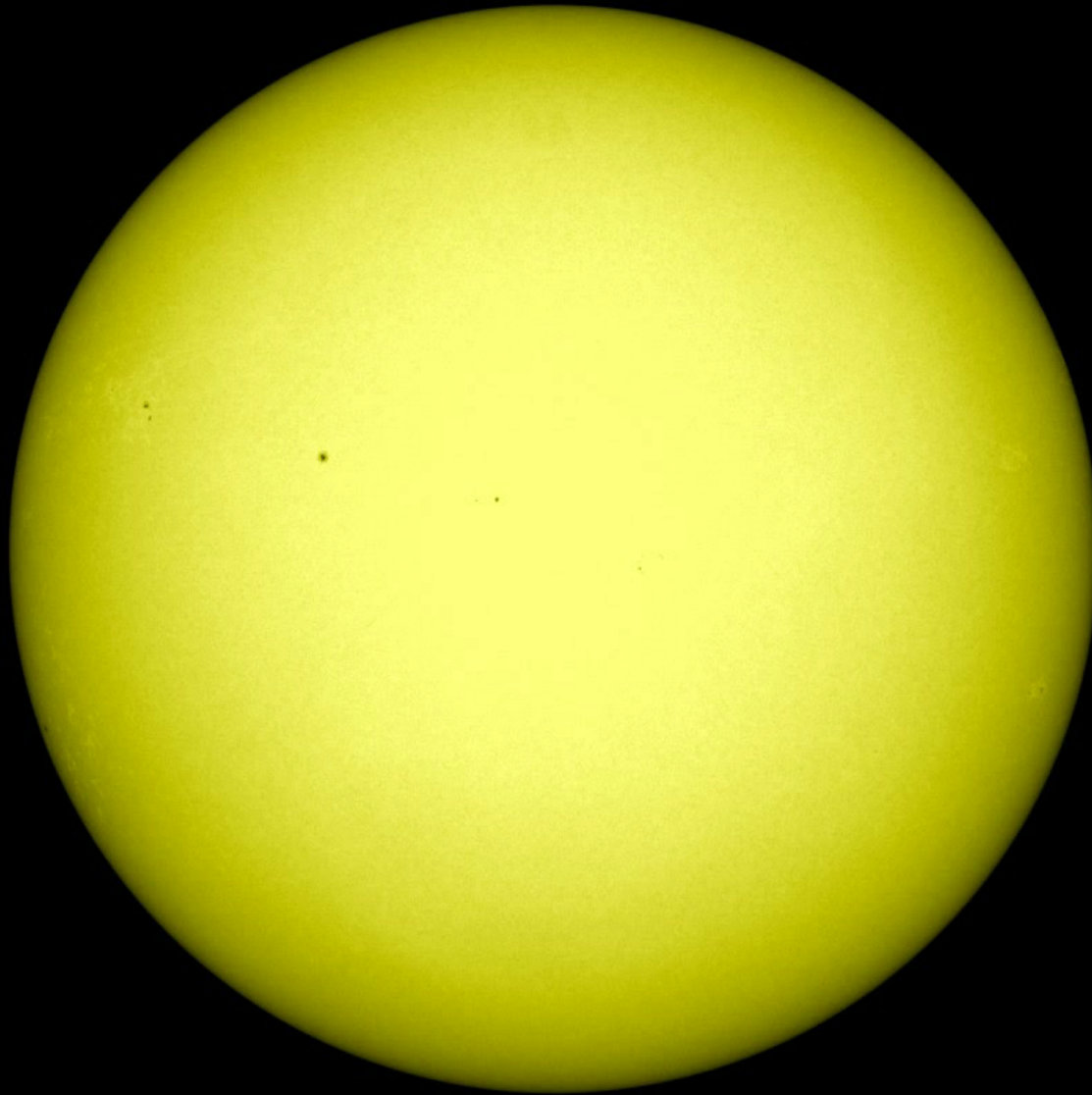
ASTR 8030/3600 Stellar Astrophysics





SDO 4500 Angstroms: photosphere

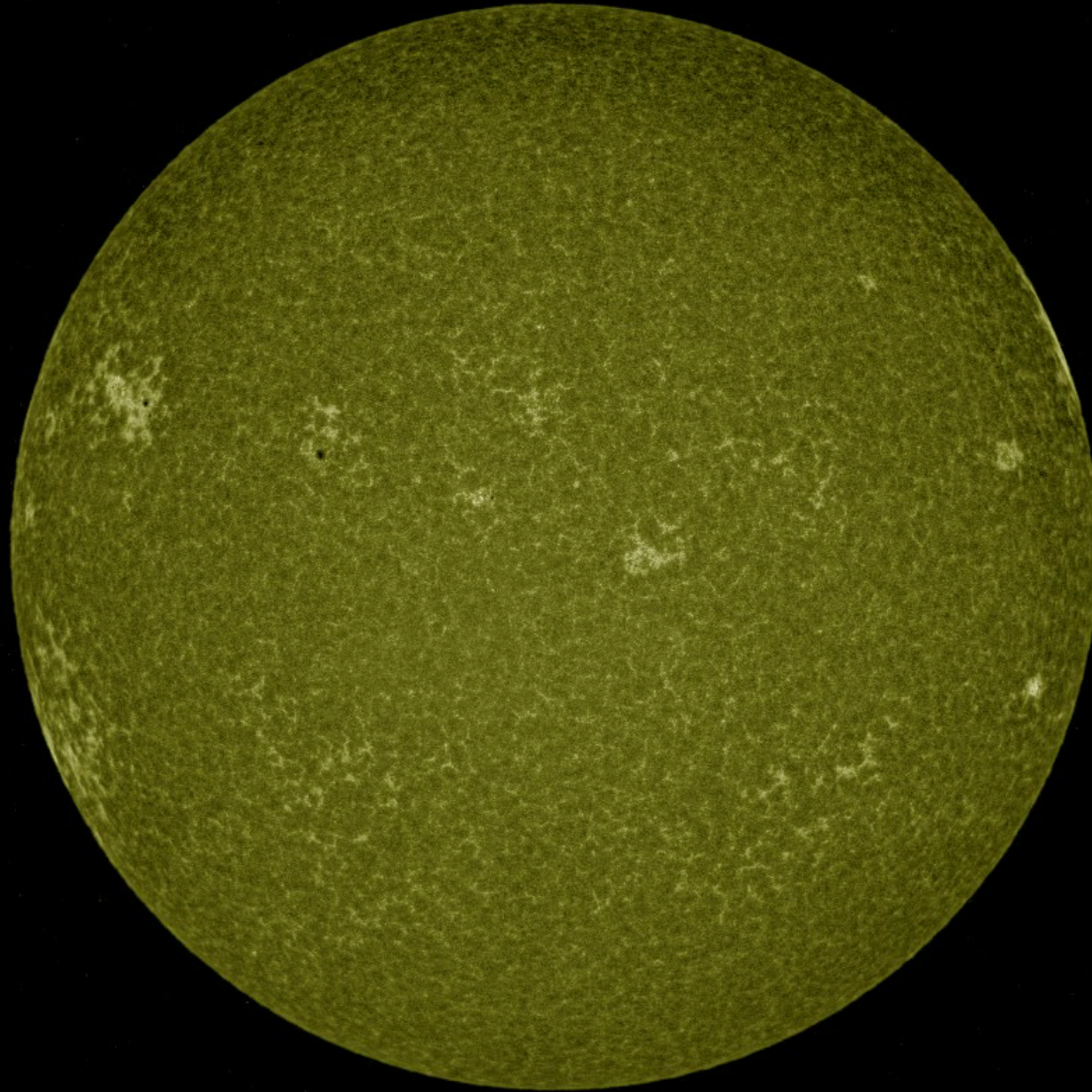
T~5000K





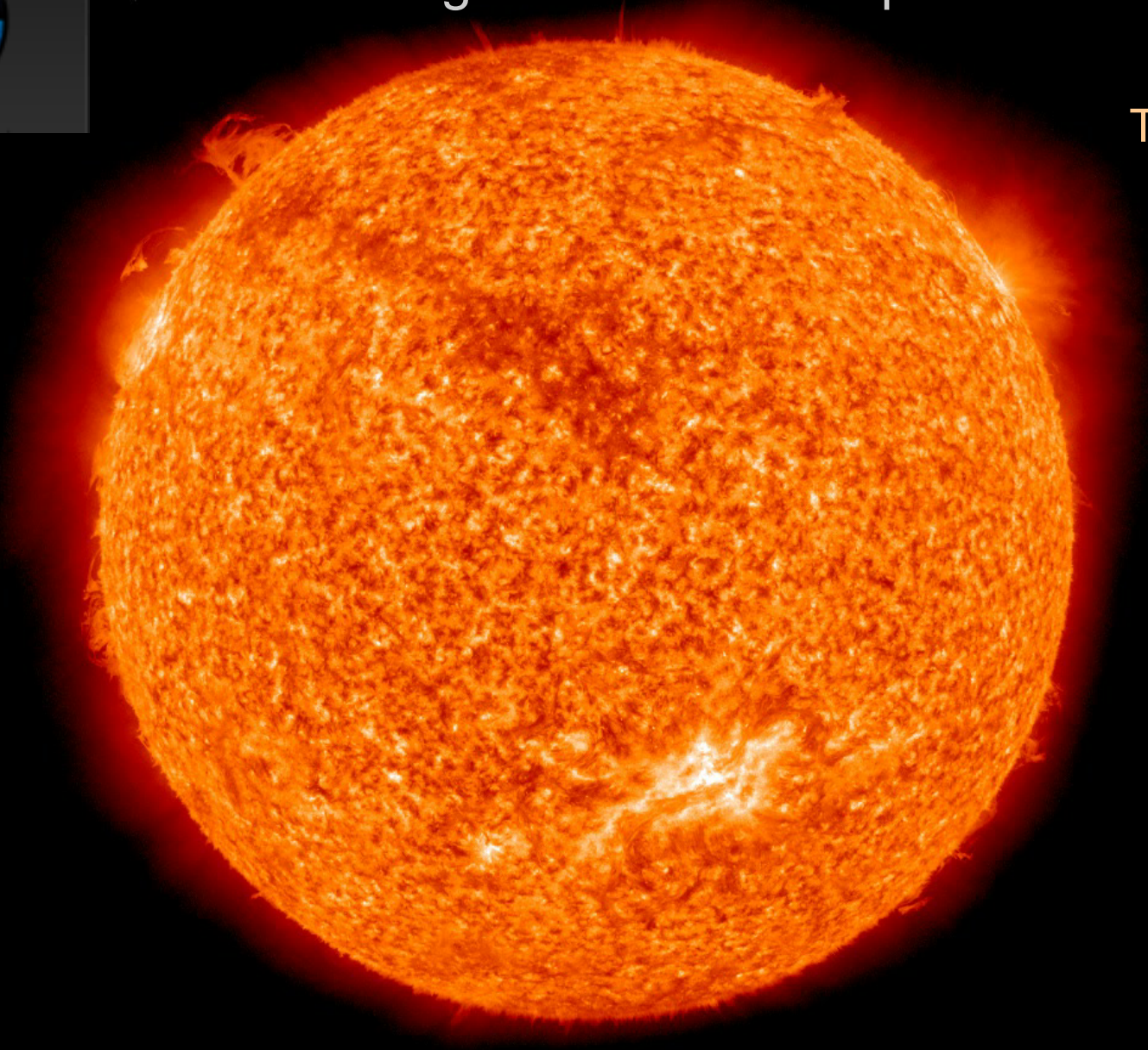
SDO 1600 Angstroms: upper photosphere

$T \sim 5 \times 10^4 \text{K}$





SDO 304 Angstroms: chromosphere

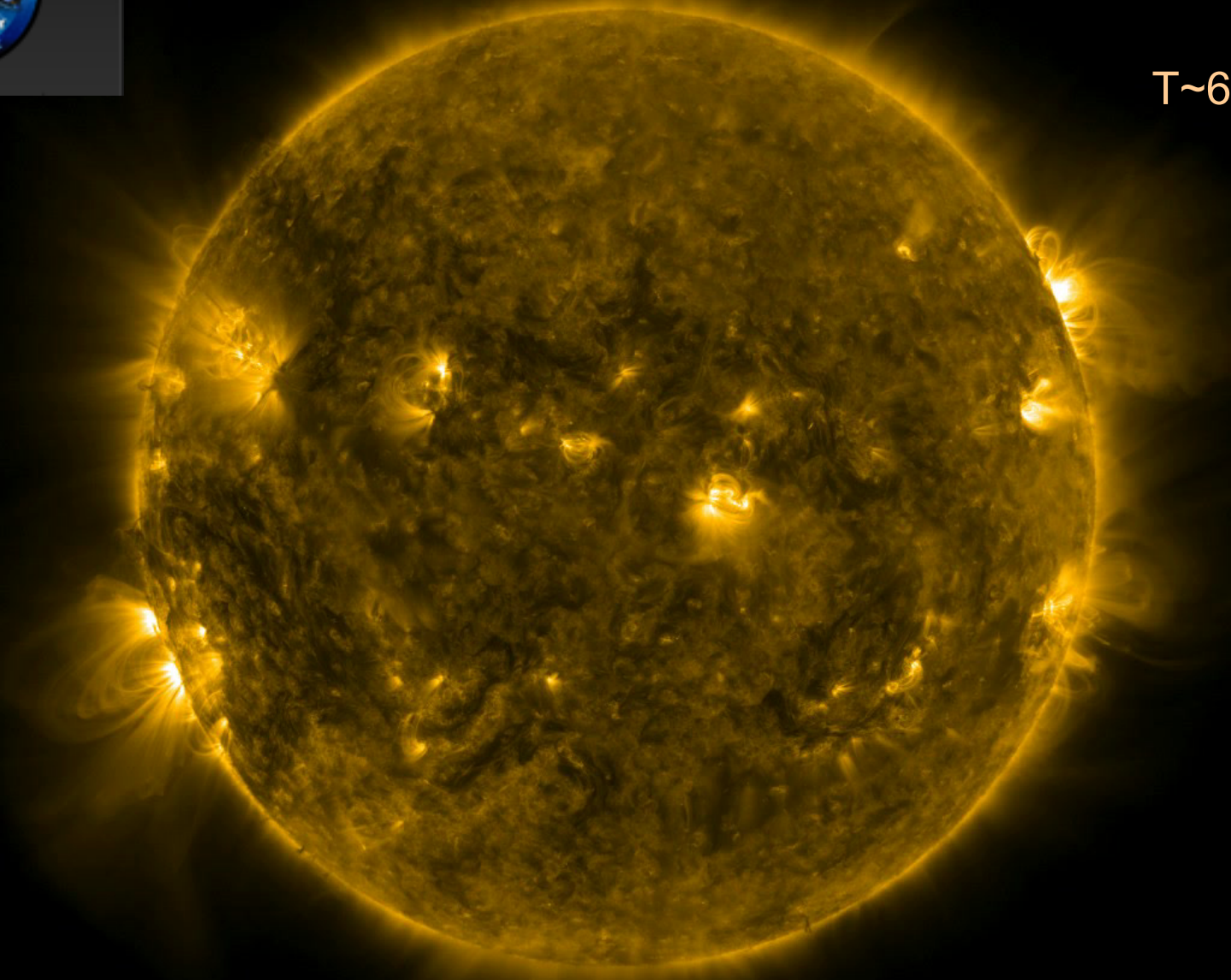


$T \sim 10^5 \text{K}$



SDO 171 Angstroms: quiet corona

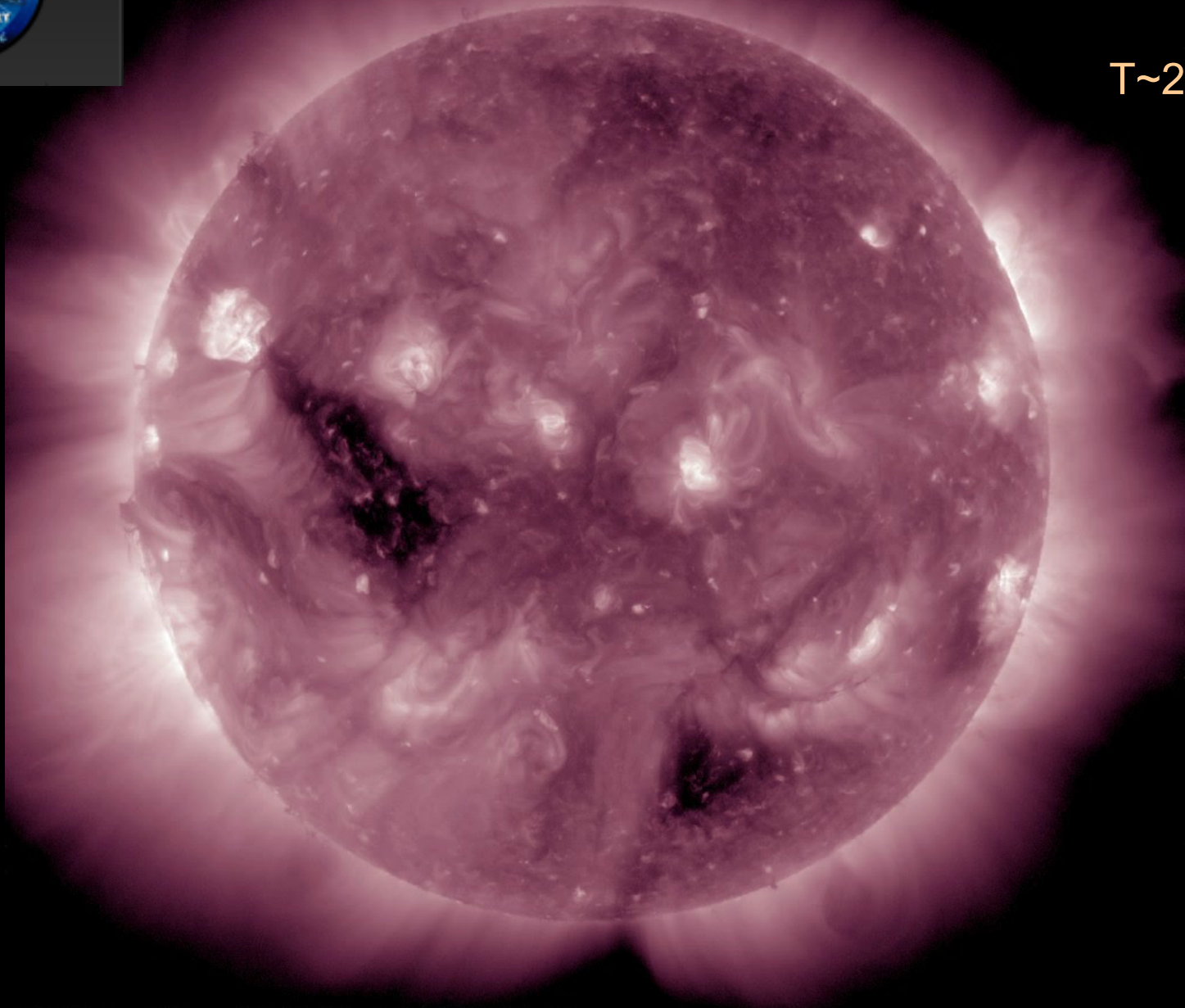
$T \sim 6 \times 10^5 \text{K}$





SDO 211 Angstroms: active corona

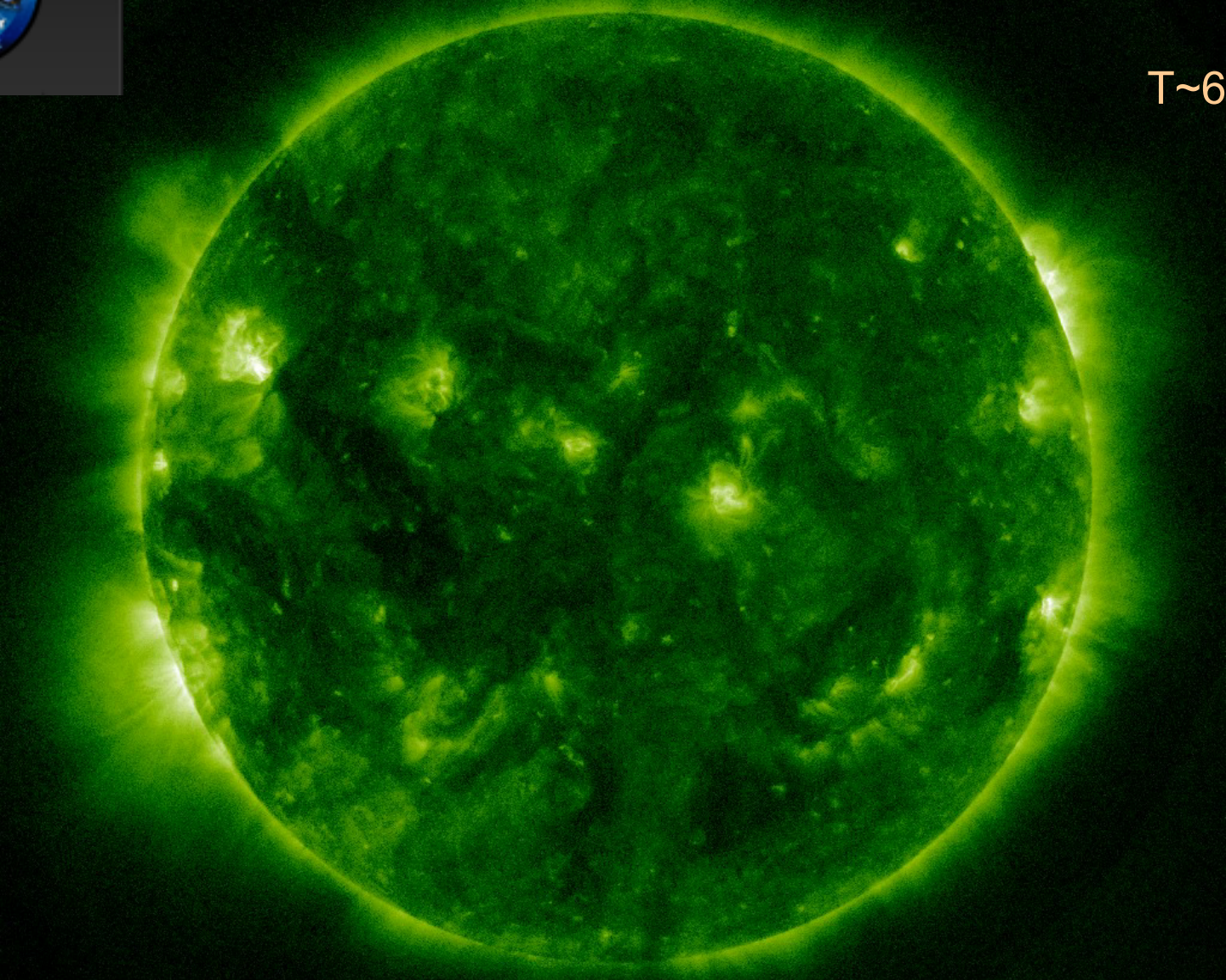
$T \sim 2 \times 10^6 \text{K}$





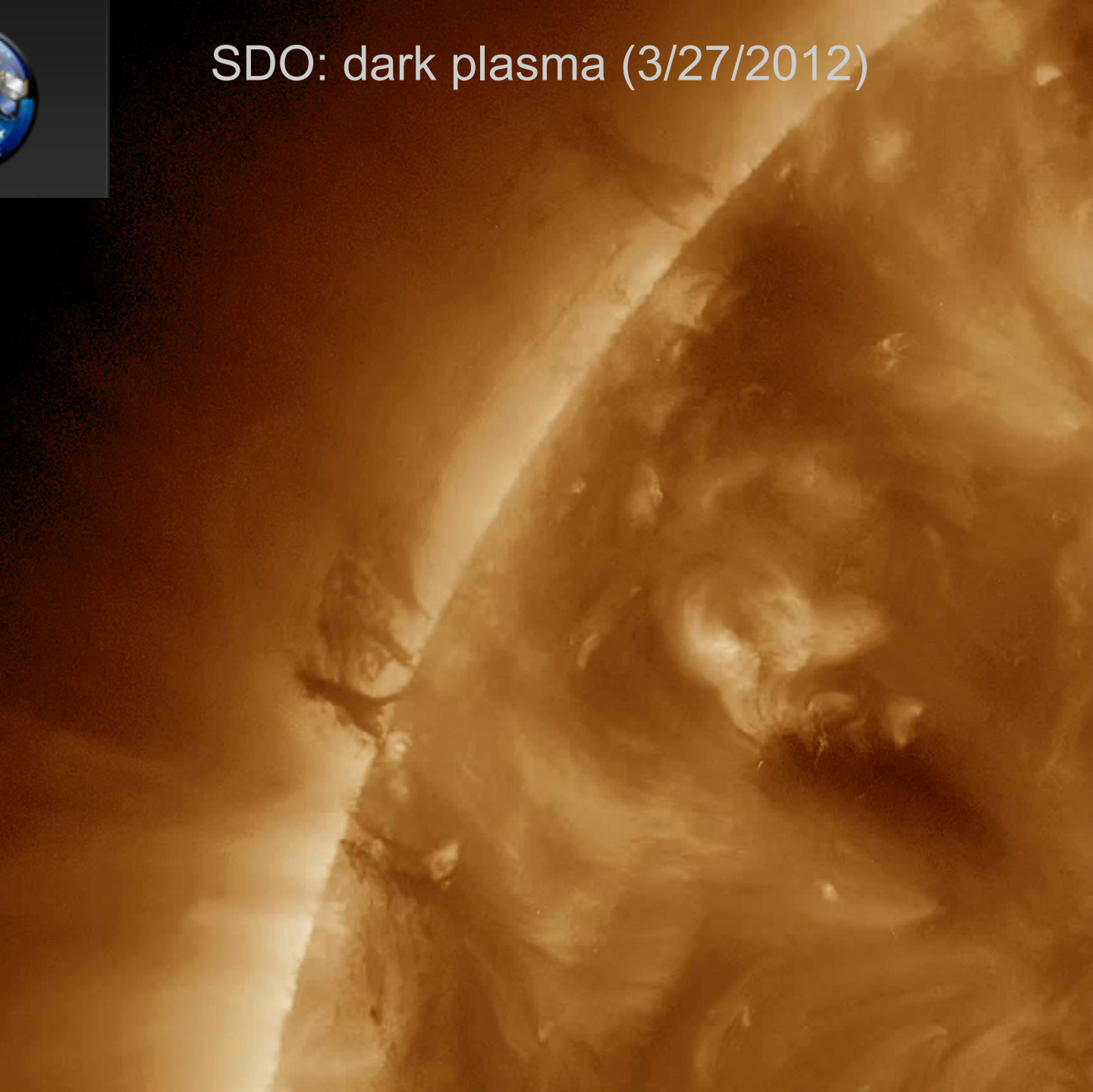
SDO 94 Angstroms: flaring regions

$T \sim 6 \times 10^6 \text{K}$



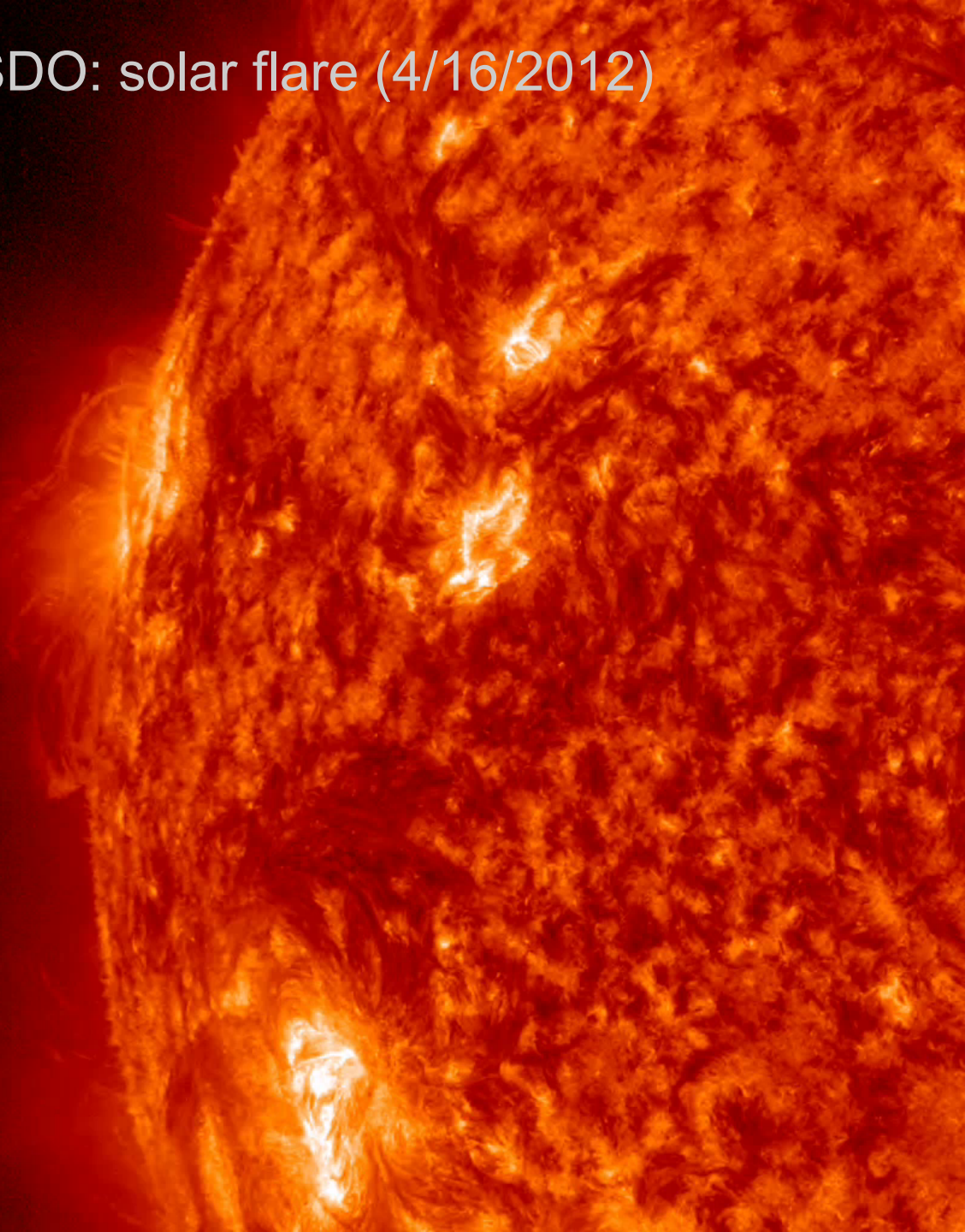


SDO: dark plasma (3/27/2012)



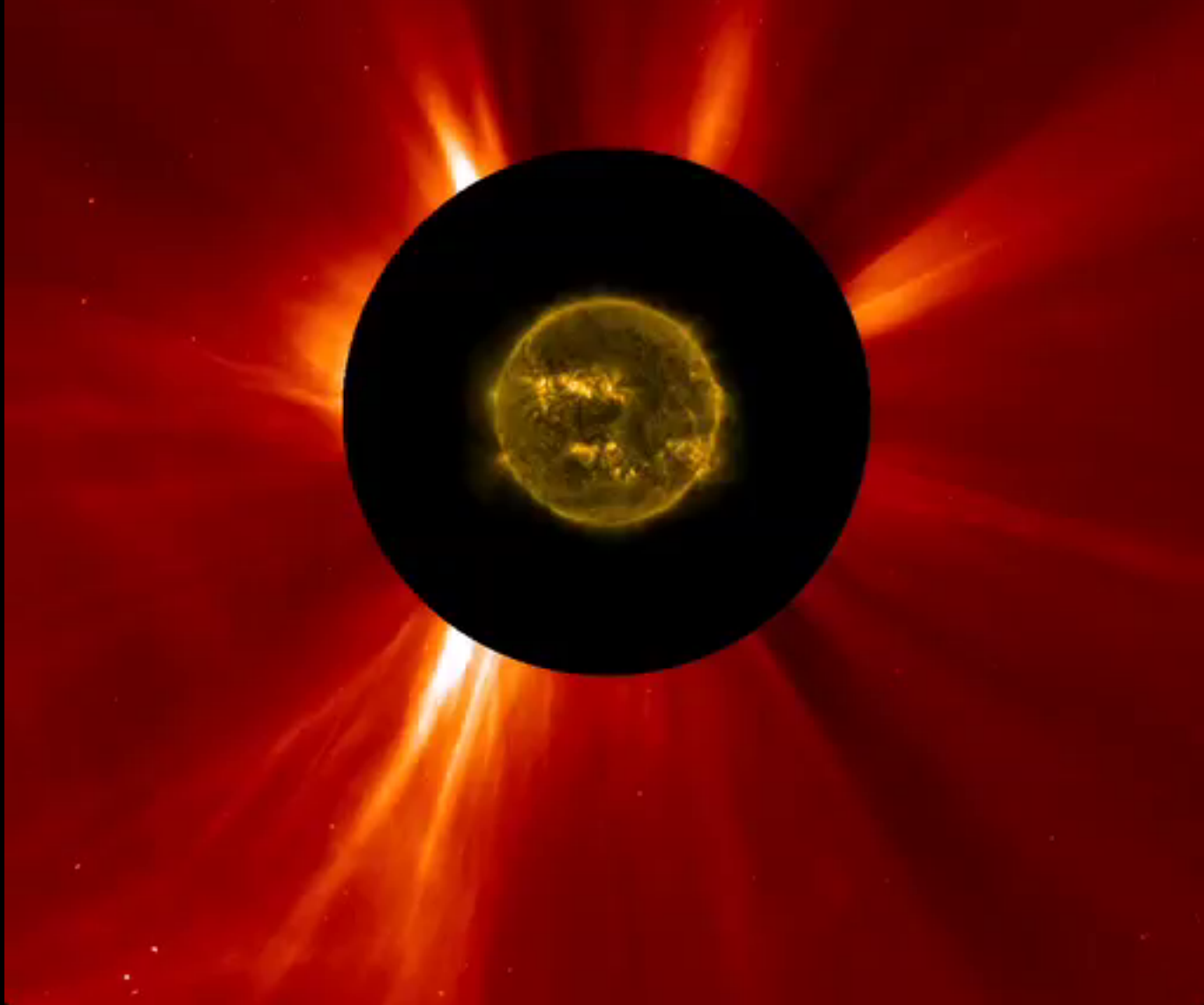


SDO: solar flare (4/16/2012)





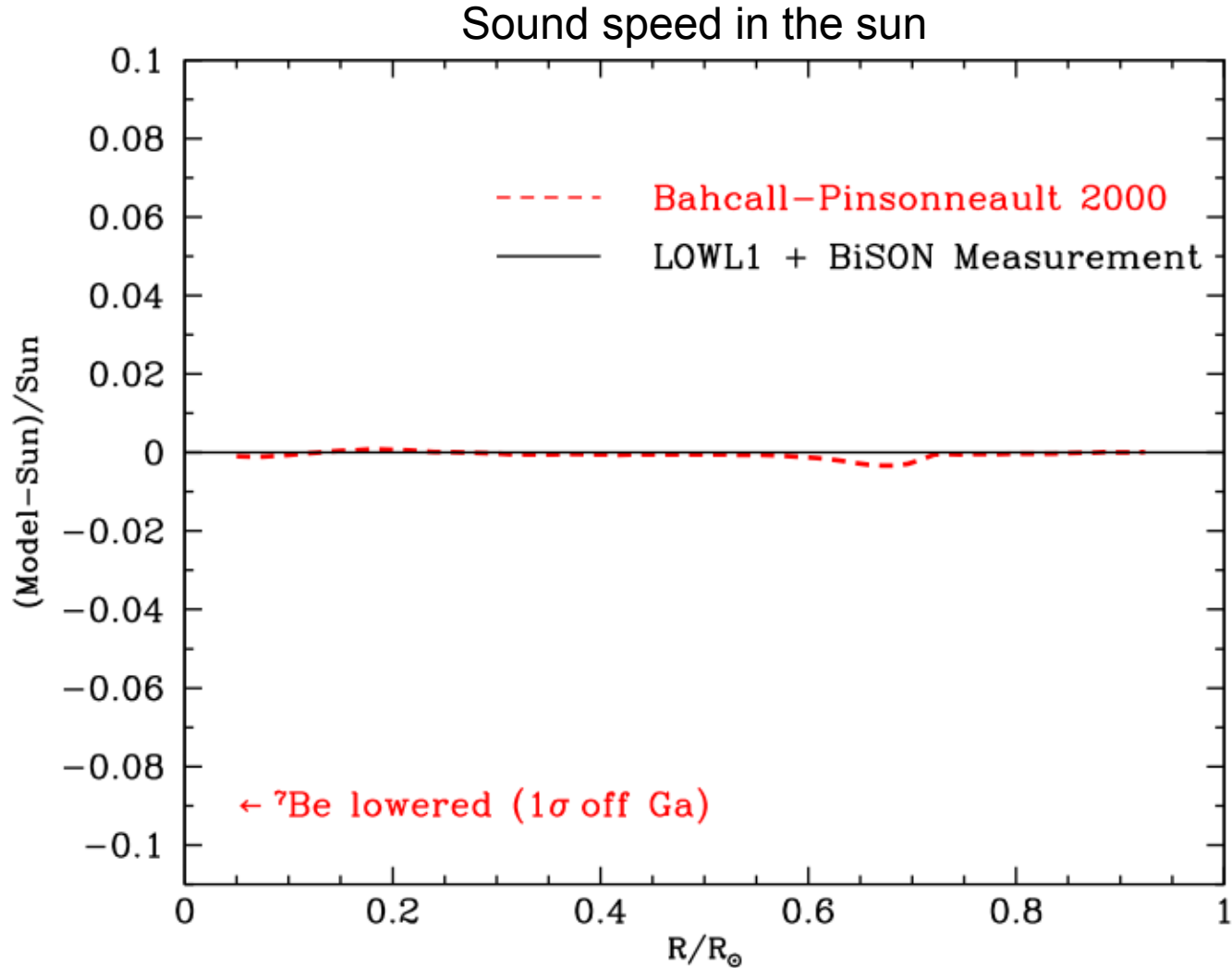
SDO: coronal mass ejection (7/2/2012)



Aims of the course

- Introduce the equations needed to model the internal structure of stars.
- Overview of how basic stellar properties are observationally measured.
- Study the microphysics relevant for stars: the equation of state, the opacity, nuclear reactions.
- Examine the properties of simple models for stars and consider how real models are computed.
- Survey (mostly qualitatively) how stars evolve, and the endpoints of stellar evolution.

Stars are relatively simple physical systems

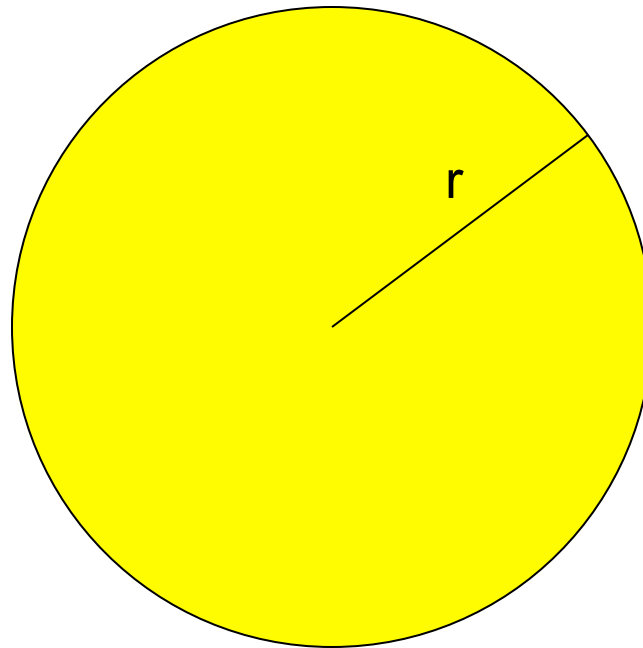


Problem of Stellar Structure

We want to determine the structure (density, temperature, energy output, pressure as a function of radius) of an isolated mass M of gas with a given composition (e.g., H, He, etc.)

Known:

Mass
+
Composition



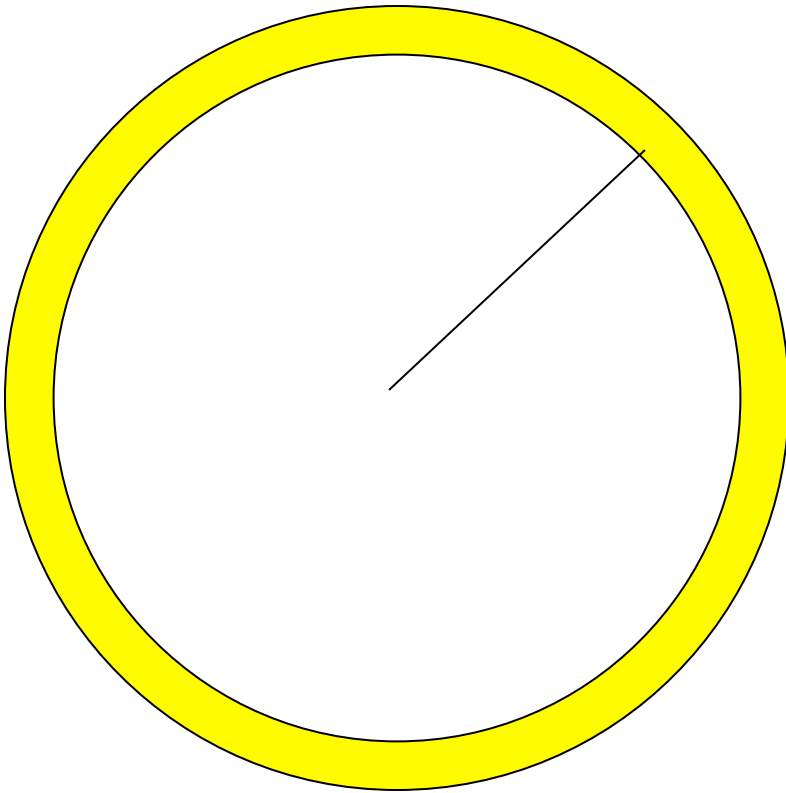
Unknown:

Density
Temperature
Energy
Pressure

Simplifying assumptions

- 1. No rotation** → spherical symmetry ✓
For sun: rotation period at surface ~ 1 month
orbital period at surface ~ few hours
- 2. No magnetic fields** ✓
For sun: magnetic field ~ 5G, ~ 1KG in sunspots
equipartition field ~ 100 MG
Some neutron stars have a large fraction of their energy in B fields
- 3. Static** ✓
For sun: convection, but no large scale variability
Not valid for forming stars, pulsating stars and dying stars.
- 4. Newtonian gravity** ✓
For sun: escape velocity ~ 600 km/s $\ll c$
Not true for neutron stars

The Variables of Stellar Structure

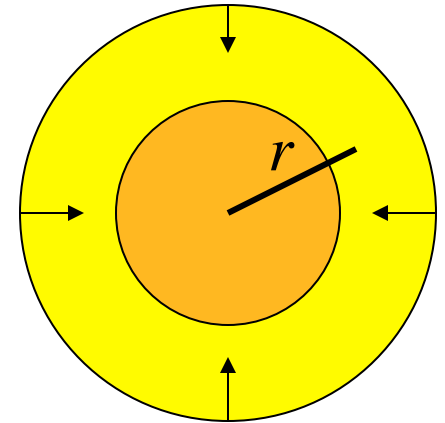
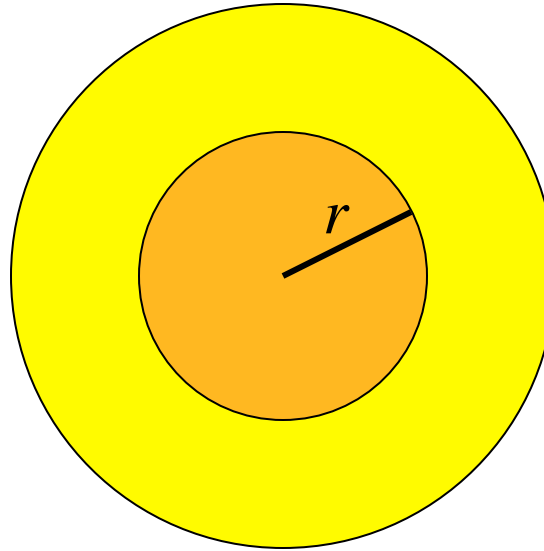


r	radius
$m(r)$	enclosed mass
$\rho(r)$	mass density
$P(r)$	Pressure
$g(r)$	gravity
$T(r)$	Temperature
$L_r(r)$	Luminosity flow
$X_i(r)$	Composition

Eulerian vs. Lagrangian

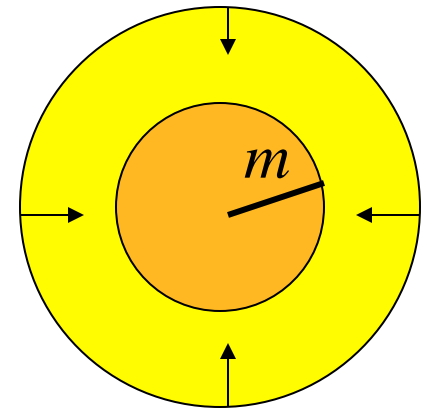
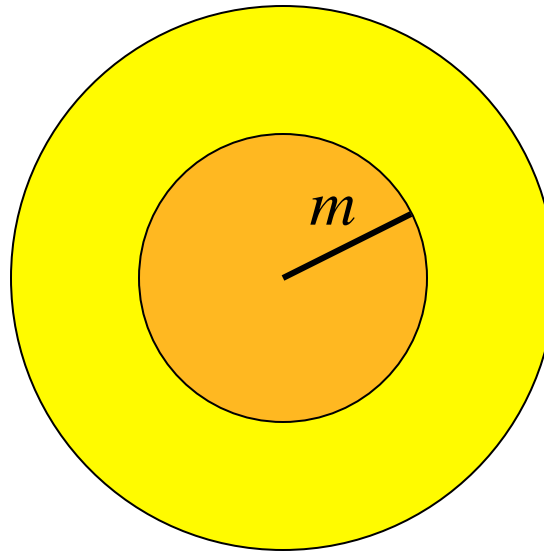
Eulerian

Characterize quantities as a function of radius.



Lagrangian

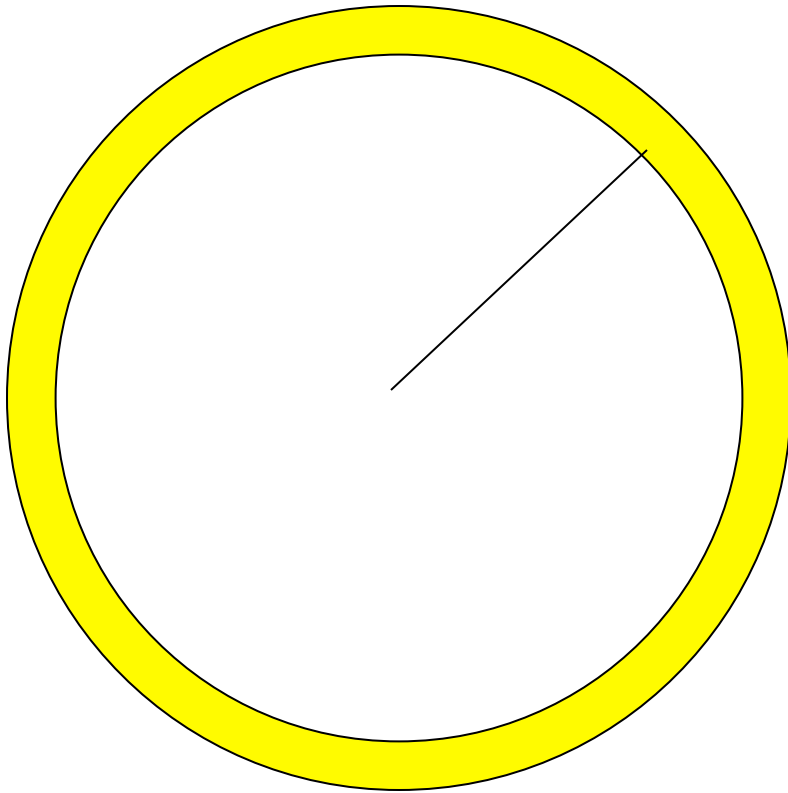
Characterize quantities as a function of enclosed mass.



time A

time B

The Variables of Stellar Structure



m enclosed mass

$r(m)$ radius

$\rho(m)$ mass density

$P(m)$ Pressure

$g(m)$ gravity

$T(m)$ Temperature

$L_r(m)$ Luminosity flow

$X_i(m)$ Composition

Values for the Sun

$$\text{Mass} \equiv 1M_{\odot} \approx 2 \times 10^{33} \text{ g}$$

$$\text{Radius} \equiv 1R_{\odot} \approx 7 \times 10^{10} \text{ cm}$$

$$\text{Luminosity} \equiv 1L_{\odot} \approx 4 \times 10^{33} \text{ erg/s}$$

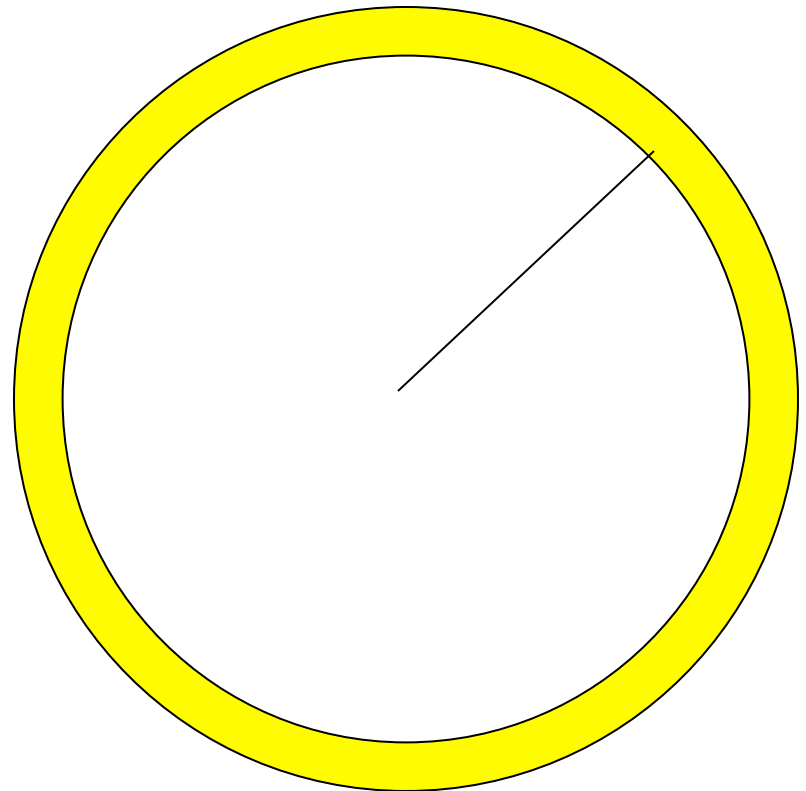
$$\text{Surface Temperature} \approx 5,800 \text{ K}$$

$$\text{Composition} \approx 70\% \text{ H} \quad 28\% \text{ He} \quad 2\% \text{ metals}$$

The Equations of Stellar Structure

mass in shell = density \times volume of shell

$$dm = \rho \times 4\pi r^2 dr$$



The Equations of Stellar Structure

1. Mass Conservation

$$\frac{dm}{dr} = 4\pi r^2 \rho$$

eulerian

$$\frac{dr}{dm} = \frac{1}{4\pi r^2 \rho}$$

lagrangian

The Equations of Stellar Structure

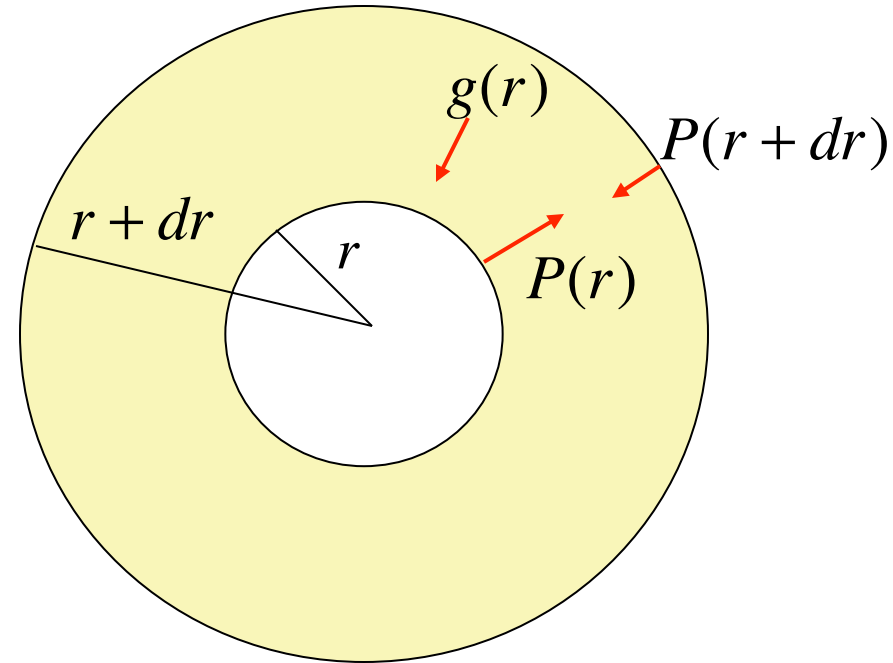
$$F = ma$$

Force pushing outwards:

$$P(r)4\pi r^2$$

Force pushing inwards:

$$-P(r + dr)4\pi r^2 - \frac{Gm}{r^2} dm$$



$$P(r)4\pi r^2 - P(r + dr)4\pi r^2 - \frac{Gm}{r^2} dm = \ddot{r} dm$$

$$\rightarrow -dP4\pi r^2 - \frac{Gm}{r^2} dm = \ddot{r} dm \quad \rightarrow -\frac{dP}{dm} 4\pi r^2 - \frac{Gm}{r^2} = \ddot{r} = 0$$

The Equations of Stellar Structure

2. Hydrostatic Equilibrium

$$\frac{dP}{dr} = -\rho \frac{Gm}{r^2}$$

eulerian

$$\frac{dP}{dm} = -\frac{Gm}{4\pi r^4}$$

lagrangian

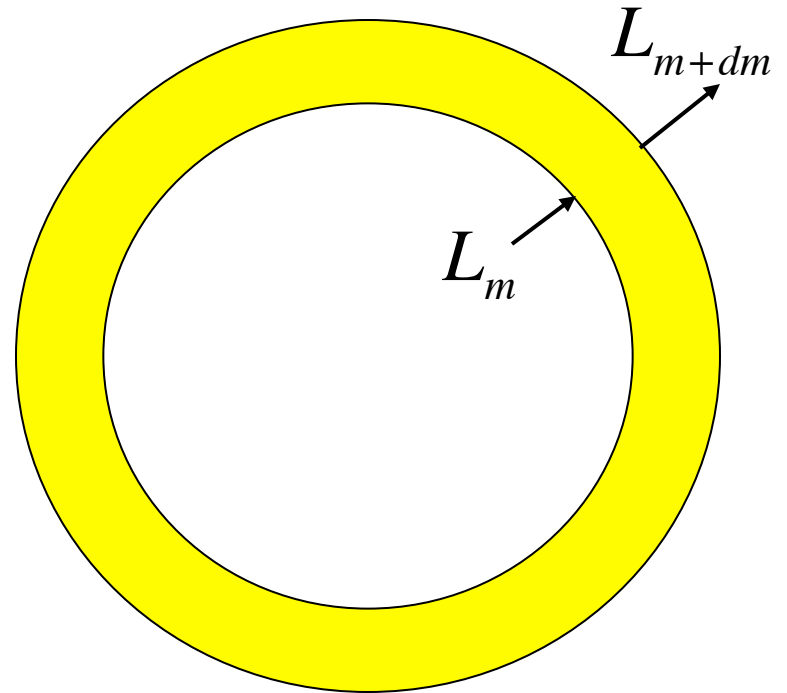
Need equation of state: $P = f(\rho, T, X_i)$

e.g., for an ideal gas: $P = \frac{R}{\mu} \rho T$

The Equations of Stellar Structure

ϵ = energy generation rate per unit mass (erg/s/g)

change in $L_m = \epsilon \times dm$



The Equations of Stellar Structure

3. Energy Generation

$$\frac{dL_r}{dr} = 4\pi r^2 \rho \epsilon$$

eulerian

$$\frac{dL_m}{dm} = \epsilon$$

lagrangian

Need nuclear physics: $\epsilon = f(\rho, T, X_i)$

e.g., for the proton-proton chain: $\epsilon \approx \epsilon_0 \rho T^4$

The Equations of Stellar Structure

Energy flow:

Depends on opacity κ (area/mass)

- radiation

low opacity

- convection

high opacity

- conduction

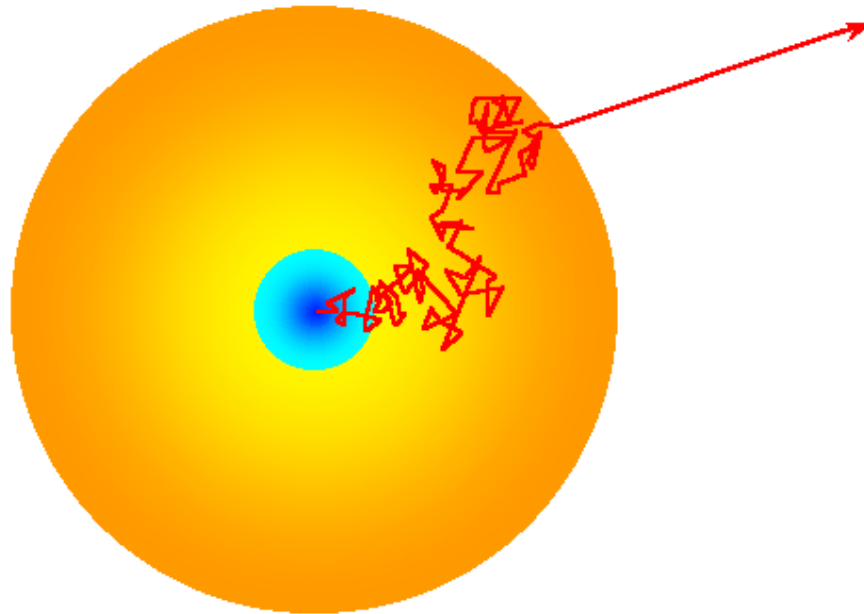
unimportant

For radiation:

radiation pressure decreases outward →

photons have net movement outward in their random walk.

Random walk of photon through the sun.



Straight path: 2.3 seconds
Random walk: 30,000 years

The Equations of Stellar Structure

4. Energy Flow (radiation)

$$\frac{dT}{dr} = - \frac{3\kappa\rho L_r}{16\pi a c r^2 T^3}$$

eulerian

$$\frac{dT}{dm} = - \frac{3\kappa L_m}{64\pi^2 a c r^4 T^3}$$

lagrangian

Need opacity: $\kappa = f(\rho, T, X_i)$

e.g., for electron Thomson scattering: $\kappa \approx \kappa_0$

The Equations of Stellar Structure

4. Energy Flow (convection)

$$\frac{dT}{dr} = - \left(1 - \frac{1}{\Gamma_2} \right) \frac{Gm\rho T}{r^2 P}$$

eulerian

$$\frac{dT}{dm} = - \left(1 - \frac{1}{\Gamma_2} \right) \frac{GmT}{4\pi r^4 P}$$

lagrangian

Need adiabatic index: $\Gamma_2 = f(\rho, T, X_i)$

e.g., for a simple ideal gas: $\Gamma_2 = 5/3$

Solving the Equations

Need microphysics: $P(\rho, T, X_i)$, $\varepsilon(\rho, T, X_i)$, $\kappa(\rho, T, X_i)$

$$\frac{dr}{dm} = \frac{1}{4\pi r^2 \rho}$$

$$\frac{dP}{dm} = -\frac{Gm}{4\pi r^4}$$

$$\frac{dL_m}{dm} = \varepsilon$$

$$\frac{dT}{dm} = -\frac{3\kappa L_m}{64\pi^2 a c r^4 T^3}$$

Solve in 1 dimension from center to surface

Boundary Conditions

Center ($r = 0$): $m = 0$, $L_r = 0$, $P = P_c$, $T = T_c$

Surface ($r = R$): $m = M$, $L_r = L$, $P \sim 0$, $T \sim 0$

Lagrangian

Center ($m = 0$): $r = 0$, $L_r = 0$, $P = ?$, $T = ?$

Surface ($m = M$): $r = ?$, $L_r = ?$, $P \sim 0$, $T \sim 0$

Boundary conditions are incomplete at each end

Complications

1. Stars are luminous

radiate away energy → must change in time
chemical composition is changing. Must account for dX_i/dt

2. Convection is very complicated and important

3. Convection can change chemical composition

4. Opacities are hard to calculate. And they matter!

Homology Relations

To calculate the luminosity or radius of a star of mass M , we must solve these differential equations. However, we can get approximate scaling relations by using *Homology*.

Assume that each differential or local quantity simply scales with the global value of that quantity.

For example,
$$\frac{dr}{dm} \sim \frac{R}{M}$$

So,
$$\frac{dr}{dm} = \frac{1}{4\pi r^2 \rho}$$
 becomes
$$\frac{R}{M} \sim \frac{1}{R^2 \bar{\rho}}$$

Homology Relations

If we also know how pressure, opacity, and nuclear generation rate scale with density and temperature, we can solve all these equations to get scaling relations. Then we can turn these scaling relations into actual equations by normalizing to the Sun.

For example, suppose we find that $L \sim M^\alpha$

This means that $L = \text{const} \times M^\alpha$

For the Sun, this is $L_\odot = \text{const} \times M_\odot^\alpha$

Dividing the two equations we get

$$\frac{L}{L_\odot} = \left(\frac{M}{M_\odot} \right)^\alpha$$