

Adiabatic Exponents

Q : specific heat (erg g⁻¹)

V_ρ : specific volume (cm³g⁻¹) = $\frac{1}{\rho}$

E : specific internal energy (erg g⁻¹)

First law of thermodynamics

$$dQ = dE + PdV_\rho = dE + Pd\left(\frac{1}{\rho}\right) = dE - \frac{P}{\rho^2}d\rho$$

$$\Gamma_1 = \left(\frac{\partial \ln P}{\partial \ln \rho}\right)_{\text{ad}} = -\left(\frac{\partial \ln P}{\partial \ln V_\rho}\right)_{\text{ad}}$$

$$\frac{\Gamma_2}{\Gamma_2 - 1} = \left(\frac{\partial \ln P}{\partial \ln T}\right)_{\text{ad}} = \frac{1}{\nabla_{\text{ad}}}$$

$$\Gamma_3 - 1 = \left(\frac{\partial \ln T}{\partial \ln \rho}\right)_{\text{ad}} = -\left(\frac{\partial \ln T}{\partial \ln V_\rho}\right)_{\text{ad}}$$

$$\nabla_{\text{ad}} = \left(\frac{\partial \ln T}{\partial \ln \rho}\right)_{\text{ad}} \left(\frac{\partial \ln \rho}{\partial \ln P}\right)_{\text{ad}}$$

$$= \frac{\Gamma_3 - 1}{\Gamma_1}$$

$$= \frac{\Gamma_2 - 1}{\Gamma_2}$$

Adiabatic Exponents

- Γ_1 describes how pressure responds to compression
(relevant for dynamical processes like pulsations)
- Γ_2 describes how temperature responds to changes in pressure
(relevant for determining whether convection takes place)
- Γ_3 describes how temperature responds to compression

Ideal Gas: $\Gamma_1 = \Gamma_2 = \Gamma_3 = \frac{5}{3}$

Radiation: $\Gamma_1 = \Gamma_2 = \Gamma_3 = \frac{4}{3}$

Γ_3 is equivalent to the γ in the γ -law equation of state

$$P = (\gamma - 1)\rho E$$

For mixtures of gas and radiation, as well as for cases where chemical reaction happen, the adiabatic exponents can differ from each other.

Adiabatic Exponents: Methodology

$$P = f(T, V_\rho)$$

$$E = g(T, V_\rho)$$

$$\text{Adiabatic change} \\ dQ = 0$$

$$dQ = dE + PdV_\rho = 0$$

$$= \left(\frac{\partial E}{\partial T} \right)_{V_\rho} dT + \left(\frac{\partial E}{\partial V_\rho} \right)_T dV_\rho + PdV_\rho = 0$$

$$dP = \left(\frac{\partial P}{\partial T} \right)_{V_\rho} dT + \left(\frac{\partial P}{\partial V_\rho} \right)_T dV_\rho = \left(\frac{\partial P}{\partial T} \right)_{ad} dT = \left(\frac{\partial P}{\partial V_\rho} \right)_{ad} dV_\rho$$

$$= \left(\frac{\Gamma_2}{\Gamma_2 - 1} \right) dT = (-\Gamma_1) dV_\rho$$

Adiabatic Exponents: Ideal gas + Radiation

Pressure

$$P = P_{\text{gas}} + P_{\text{rad}} = \frac{N_A k}{\mu} \rho T + \frac{1}{3} a T^4$$

$$= \frac{N_A k}{\mu} \frac{T}{V_\rho} + \frac{1}{3} a T^4$$

Specific Internal Energy

$$E = E_{\text{gas}} + E_{\text{rad}} = \frac{3}{2} \frac{N_A k}{\mu} T + a \frac{T^4}{\rho}$$

$$= \frac{3}{2} \frac{N_A k}{\mu} T + a T^4 V_\rho$$

Adiabatic Exponents: Ideal gas + Radiation

$$P = \frac{N_A k T}{\mu V_\rho} + \frac{1}{3} a T^4$$

$$E = \frac{3 N_A k}{2 \mu} T + a T^4 V_\rho$$

$$\text{Adiabatic change} \\ dQ = 0$$

$$\begin{aligned} dQ &= \left(\frac{\partial E}{\partial T} \right)_{V_\rho} dT + \left(\frac{\partial E}{\partial V_\rho} \right)_T dV_\rho + P dV_\rho = 0 \\ &= \left(\frac{3 N_A k}{2 \mu} + 4 a T^3 V_\rho \right) dT + (a T^4) dV_\rho + \left(\frac{N_A k T}{\mu V_\rho} + \frac{1}{3} a T^4 \right) dV_\rho = 0 \\ &= \left(\frac{3 N_A k}{2 \mu} + 4 a T^3 V_\rho \right) dT + \left(\frac{N_A k T}{\mu V_\rho} + \frac{4}{3} a T^4 \right) dV_\rho = 0 \\ &= \left(\frac{3 N_A k T}{2 \mu V_\rho} + 4 a T^4 \right) \frac{V_\rho}{T} dT + \left(\frac{N_A k T}{\mu V_\rho} + \frac{4}{3} a T^4 \right) dV_\rho = 0 \end{aligned}$$

$$= \left(\frac{3}{2} P_{\text{gas}} + 12 P_{\text{rad}} \right) \frac{dT}{T} + \left(P_{\text{gas}} + 4 P_{\text{rad}} \right) \frac{dV_\rho}{V_\rho} = 0$$

Adiabatic Exponents: Ideal gas + Radiation

$$P = \frac{N_A k T}{\mu V_\rho} + \frac{1}{3} a T^4$$

$$E = \frac{3 N_A k}{2 \mu} T + a T^4 V_\rho$$

Equation
of state

$$dP = \left(\frac{\partial P}{\partial T} \right)_{V_\rho} dT + \left(\frac{\partial P}{\partial V_\rho} \right)_T dV_\rho = \left(\frac{N_A k}{\mu} \frac{1}{V_\rho} + \frac{4}{3} a T^3 \right) dT + \left(-\frac{N_A k T}{\mu V_\rho^2} \right) dV_\rho$$

$$= \left(\frac{N_A k T}{\mu V_\rho} + \frac{4}{3} a T^4 \right) \frac{1}{T} dT + \left(-\frac{N_A k T}{\mu V_\rho} \right) \frac{1}{V_\rho} dV_\rho$$

$$= (P_{\text{gas}} + 4P_{\text{rad}}) \frac{dT}{T} - P_{\text{gas}} \frac{dV_\rho}{V_\rho} = dP = \left(\frac{\partial P}{\partial T} \right)_{\text{ad}} dT = \frac{P}{T} \left(\frac{\partial \ln P}{\partial \ln T} \right)_{\text{ad}} dT$$

$$\rightarrow (P_{\text{gas}} + 4P_{\text{rad}}) \frac{dT}{T} - P_{\text{gas}} \frac{dV_\rho}{V_\rho} - \frac{\Gamma_2}{\Gamma_2 - 1} (P_{\text{gas}} + P_{\text{rad}}) \frac{dT}{T} = 0$$

$$= \left(P_{\text{gas}} + 4P_{\text{rad}} - \frac{\Gamma_2}{\Gamma_2 - 1} (P_{\text{gas}} + P_{\text{rad}}) \right) \frac{dT}{T} - P_{\text{gas}} \frac{dV_\rho}{V_\rho} = 0$$

Adiabatic Exponents: Ideal gas + Radiation

1st Law of Thermodynamics

$$\left(\frac{3}{2} P_{\text{gas}} + 12 P_{\text{rad}} \right) \frac{dT}{T} + (P_{\text{gas}} + 4 P_{\text{rad}}) \frac{dV_{\rho}}{V_{\rho}} = 0$$

Equation of State

$$\left(P_{\text{gas}} + 4 P_{\text{rad}} - \frac{\Gamma_2}{\Gamma_2 - 1} (P_{\text{gas}} + P_{\text{rad}}) \right) \frac{dT}{T} - P_{\text{gas}} \frac{dV_{\rho}}{V_{\rho}} = 0$$

$$\left. \begin{array}{l} Ax + By = 0 \\ Cx + Dy = 0 \end{array} \right\} \rightarrow \frac{A}{C} = \frac{B}{D}$$

$$\frac{P_{\text{gas}} + 4 P_{\text{rad}} - \frac{\Gamma_2}{\Gamma_2 - 1} (P_{\text{gas}} + P_{\text{rad}})}{\frac{3}{2} P_{\text{gas}} + 12 P_{\text{rad}}} = - \frac{P_{\text{gas}}}{P_{\text{gas}} + 4 P_{\text{rad}}}$$

$$\begin{array}{l} P_{\text{gas}} = \beta P \\ P_{\text{rad}} = (1 - \beta) P \end{array}$$

$$\frac{\cancel{\beta P} + 4(1 - \cancel{\beta})\cancel{P} - \frac{\Gamma_2}{\Gamma_2 - 1} \cancel{P}}{\frac{3}{2} \cancel{\beta P} + 12(1 - \cancel{\beta})\cancel{P}} = - \frac{\cancel{\beta P}}{\cancel{\beta P} + 4(1 - \cancel{\beta})\cancel{P}}$$

Adiabatic Exponents: Ideal gas + Radiation

$$\frac{\beta + 4(1 - \beta) - \frac{\Gamma_2}{\Gamma_2 - 1}}{\frac{3}{2}\beta + 12(1 - \beta)} = -\frac{\beta}{\beta + 4(1 - \beta)} \rightarrow \frac{4 - 3\beta - \frac{\Gamma_2}{\Gamma_2 - 1}}{12 - \frac{21}{2}\beta} = -\frac{\beta}{4 - 3\beta}$$

$$\rightarrow \left(4 - 3\beta - \frac{\Gamma_2}{\Gamma_2 - 1}\right)(4 - 3\beta) = (-\beta)\left(12 - \frac{21}{2}\beta\right)$$

$$\rightarrow 16 - 12\beta - 12\beta + 9\beta^2 - \frac{\Gamma_2}{\Gamma_2 - 1}(4 - 3\beta) = -12\beta + \frac{21}{2}\beta^2$$

$$\rightarrow \frac{\Gamma_2}{\Gamma_2 - 1}(4 - 3\beta) = 16 - 12\beta - \frac{3}{2}\beta^2 \quad \rightarrow \frac{\Gamma_2 - 1}{\Gamma_2} = \frac{4 - 3\beta}{16 - 12\beta - \frac{3}{2}\beta^2}$$

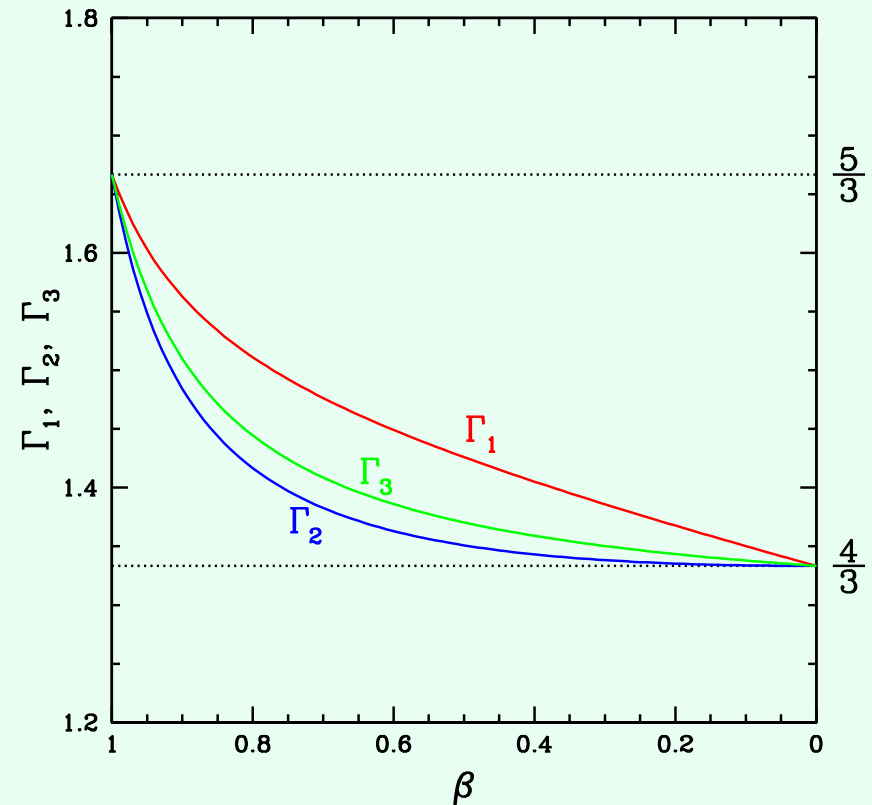
$$\rightarrow \frac{1}{\Gamma_2} = 1 - \frac{4 - 3\beta}{16 - 12\beta - \frac{3}{2}\beta^2} = \frac{16 - 12\beta - \frac{3}{2}\beta^2 - 4 + 3\beta}{16 - 12\beta - \frac{3}{2}\beta^2} = \frac{12 - 9\beta - \frac{3}{2}\beta^2}{16 - 12\beta - \frac{3}{2}\beta^2}$$

Adiabatic Exponents: Ideal gas + Radiation

$$\frac{1}{\Gamma_2} = \frac{12 - 9\beta - \frac{3}{2}\beta^2}{16 - 12\beta - \frac{3}{2}\beta^2} \rightarrow \Gamma_2 = \frac{16 - 12\beta - \frac{3}{2}\beta^2}{12 - 9\beta - \frac{3}{2}\beta^2} \rightarrow \Gamma_2 = \frac{32 - 24\beta - 3\beta^2}{24 - 18\beta - 3\beta^2}$$

Similarly, we get:

$$\Gamma_1 = \frac{32 - 24\beta - 3\beta^2}{24 - 21\beta}$$
$$\Gamma_2 = \frac{32 - 24\beta - 3\beta^2}{24 - 18\beta - 3\beta^2}$$
$$\Gamma_3 = \frac{32 - 27\beta}{24 - 21\beta}$$



Adiabatic Exponents: Ionization Zones

Thermodynamics of partially ionized gas is different because

- The number of free particles is not constant
- Ionization energy is required to increase n

Assuming a pure H ideal gas

$$\left. \begin{aligned} n^+ &= yn_I \\ n^0 &= (1-y)n_I \\ n_e &= yn_I \\ n_I &= N_A \rho \end{aligned} \right\} \frac{y^2}{1-y} = \frac{1}{n_I} \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} e^{-\chi/kT}$$
$$= \frac{A}{N_A \rho} T^{3/2} e^{-\chi/kT}$$

$$\frac{y^2}{1-y} = \frac{A}{N_A} V_\rho T^{3/2} e^{-\chi/kT}$$

Adiabatic Exponents: Ionization Zones

Pressure

$$P = (n^0 + n^+ + n_e)kT = [(1-y)n_I + yn_I + yn_I]kT$$

$$= (1+y)n_IkT = (1+y)N_A\rho kT \rightarrow$$

$$P = (1+y)N_Ak\frac{T}{V_\rho}$$

Specific Internal Energy

$$E = \frac{3}{2}(n^0 + n^+ + n_e)\frac{kT}{\rho} + \frac{n^+\chi}{\rho} = (1+y)n_I\frac{3}{2}\frac{kT}{\rho} + \frac{yn_I}{\rho}\chi$$

$$E = (1+y)\frac{3}{2}N_AkT + yN_A\chi$$

Adiabatic Exponents: Ionization Zones

$$P = (1 + y) N_A k \frac{T}{V_\rho}$$

$$E = (1 + y) \frac{3}{2} N_A k T + y N_A \chi$$

$$\text{Adiabatic change} \\ dQ = 0$$

$$\begin{aligned} dQ &= \left(\frac{\partial E}{\partial T} \right)_{V_\rho, y} dT + \left(\frac{\partial E}{\partial V_\rho} \right)_{T, y} dV_\rho + \left(\frac{\partial E}{\partial y} \right)_{T, V_\rho} dy + P dV_\rho = 0 \\ &= \left[(1 + y) \frac{3}{2} N_A k \right] dT + \left[\frac{3}{2} N_A k T + N_A \chi \right] dy + \left[(1 + y) N_A k \frac{T}{V_\rho} \right] dV_\rho = 0 \\ &= \left[(1 + y) \frac{3}{2} N_A k T \right] \frac{dT}{T} + \left[(1 + y) \frac{3}{2} N_A k T + (1 + y) N_A \chi \right] \frac{dy}{1 + y} \\ &\quad + \left[(1 + y) N_A k T \right] \frac{dV_\rho}{V_\rho} = 0 \end{aligned}$$

$$= E_{\text{ideal}} \left[\frac{dT}{T} + \frac{2}{3} \frac{dV_\rho}{V_\rho} + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) \frac{dy}{1 + y} \right] = 0$$

Adiabatic Exponents: Ionization Zones

$$P = (1 + y) N_A k \frac{T}{V_\rho}$$

$$E = (1 + y) \frac{3}{2} N_A k T + y N_A \chi$$

Equation
of State

$$dP = \left(\frac{\partial P}{\partial T} \right)_{V_\rho, y} dT + \left(\frac{\partial P}{\partial V_\rho} \right)_{T, y} dV_\rho + \left(\frac{\partial P}{\partial y} \right)_{T, V_\rho} dy$$

$$= \left[(1 + y) N_A k \frac{1}{V_\rho} \right] dT - \left[(1 + y) N_A k \frac{T}{V_\rho^2} \right] dV_\rho + \left[N_A k \frac{T}{V_\rho} \right] dy$$

$$= \left[(1 + y) N_A k \frac{T}{V_\rho} \right] \frac{dT}{T} - \left[(1 + y) N_A k \frac{T}{V_\rho} \right] \frac{dV_\rho}{V_\rho} + \left[(1 + y) N_A k \frac{T}{V_\rho} \right] \frac{dy}{1 + y}$$

$$dP = (1 + y) N_A k \frac{T}{V_\rho} \left[\frac{dT}{T} - \frac{dV_\rho}{V_\rho} + \frac{dy}{1 + y} \right] \rightarrow dP = P \left[\frac{dT}{T} - \frac{dV_\rho}{V_\rho} + \frac{dy}{1 + y} \right]$$

Adiabatic Exponents: Ionization Zones

$$f(y) = \frac{y^2}{1-y} = \frac{A}{N_A} V_\rho T^{3/2} e^{-\chi/kT}$$

Saha Equation

$$dy = \frac{dy}{df} df = \frac{dy}{df} \left[\left(\frac{\partial f}{\partial T} \right)_{V_\rho} dT + \left(\frac{\partial f}{\partial V_\rho} \right)_T dV_\rho \right]$$

$$= \frac{(1-y)^2}{2y(1-y) + y^2} \left[\frac{A}{N_A} V_\rho \left(\frac{3}{2} T^{1/2} e^{-\chi/kT} + T^{3/2} e^{-\chi/kT} \frac{\chi}{kT^2} \right) dT + \frac{A}{N_A} T^{3/2} e^{-\chi/kT} dV_\rho \right]$$

$$= \frac{(1-y)^2}{2y(1-y) + y^2} \frac{A}{N_A} V_\rho T^{3/2} e^{-\chi/kT} \left[\left(\frac{3}{2} \frac{1}{T} + \frac{\chi}{kT^2} \right) dT + \frac{dV_\rho}{V_\rho} \right]$$

$$= \frac{(1-y)^2}{2y(1-y) + y^2} \frac{y^2}{1-y} \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right]$$

Adiabatic Exponents: Ionization Zones

$$dy = \frac{(1-y)^2}{2y(1-y) + y^2} \frac{y^2}{1-y} \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right]$$

$$dy = \frac{(1-y)y}{2(1-y) + y} \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right]$$

$$dy = \frac{(1-y)y}{2-y} \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right]$$

$$\frac{dy}{1+y} = \frac{(1-y)y}{(1+y)(2-y)} \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right]$$

$D(y)$

Adiabatic Exponents: Ionization Zones

1st Law of
Thermodynamics

$$E_{\text{ideal}} \left[\frac{dT}{T} + \frac{2}{3} \frac{dV_{\rho}}{V_{\rho}} + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) \frac{dy}{1+y} \right] = 0$$

Equation
of State

$$dP = P \left[\frac{dT}{T} - \frac{dV_{\rho}}{V_{\rho}} + \frac{dy}{1+y} \right]$$

Saha

$$\frac{dy}{1+y} = D(y) \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_{\rho}}{V_{\rho}} \right]$$

Adiabatic Exponents: Ionization Zones

1st Law of
Thermodynamics

$$E_{\text{ideal}} \left[\frac{dT}{T} + \frac{2}{3} \frac{dV_{\rho}}{V_{\rho}} + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) \frac{dy}{1+y} \right] = 0$$

$$E_{\text{ideal}} \left[\frac{dT}{T} + \frac{2}{3} \frac{dV_{\rho}}{V_{\rho}} + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) D(y) \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_{\rho}}{V_{\rho}} \right] \right] = 0$$

$$E_{\text{ideal}} \left[\left[1 + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dT}{T} + \left[\frac{2}{3} + \left(1 + \frac{2}{3} \frac{\chi}{kT} \right) D(y) \right] \frac{dV_{\rho}}{V_{\rho}} \right] = 0$$

$$E_{\text{ideal}} \left[\left[1 + D(y) \frac{2}{3} \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2 \right] \frac{dT}{T} + \frac{2}{3} \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dV_{\rho}}{V_{\rho}} \right] = 0$$

Adiabatic Exponents: Ionization Zones

Equation of State

$$dP = P \left[\frac{dT}{T} - \frac{dV_\rho}{V_\rho} + \frac{dy}{1+y} \right]$$

$$dP = P \left[\frac{dT}{T} - \frac{dV_\rho}{V_\rho} + D(y) \left[\left(\frac{3}{2} + \frac{\chi}{kT} \right) \frac{dT}{T} + \frac{dV_\rho}{V_\rho} \right] \right]$$

$$dP = P \left[\left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dT}{T} + [D(y) - 1] \frac{dV_\rho}{V_\rho} \right]$$

$$dP = \left(\frac{\partial P}{\partial T} \right)_{\text{ad}} dT = \frac{P}{T} \left(\frac{\partial \ln P}{\partial \ln T} \right)_{\text{ad}} dT = \frac{P}{T} \frac{\Gamma_2}{\Gamma_2 - 1} dT$$

$$\frac{\Gamma_2}{\Gamma_2 - 1} \frac{dT}{T} = \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dT}{T} + [D(y) - 1] \frac{dV_\rho}{V_\rho}$$

$$\left[1 - \frac{\Gamma_2}{\Gamma_2 - 1} + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dT}{T} + [D(y) - 1] \frac{dV_\rho}{V_\rho} = 0$$

Adiabatic Exponents: Ionization Zones

1st Law

$$\left[1 + D(y) \frac{2}{3} \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2 \right] \frac{dT}{T} + \frac{2}{3} \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dV_\rho}{V_\rho} = 0$$

Equation
of State

$$\left[1 - \frac{\Gamma_2}{\Gamma_2 - 1} + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right] \frac{dT}{T} + [D(y) - 1] \frac{dV_\rho}{V_\rho} = 0$$

$$\frac{1 - \frac{\Gamma_2}{\Gamma_2 - 1} + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right)}{1 + D(y) \frac{2}{3} \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2} = \frac{D(y) - 1}{\frac{2}{3} \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right]}$$

$$\frac{\Gamma_2}{\Gamma_2 - 1} = 1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) - \frac{(D(y) - 1) \left[1 + D(y) \frac{2}{3} \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2 \right]}{\frac{2}{3} \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right]}$$

Adiabatic Exponents: Ionization Zones

$$\frac{\Gamma_2}{\Gamma_2 - 1} = 1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) - \frac{(D(y) - 1) \left[1 + D(y) \frac{2}{3} \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2 \right]}{\frac{2}{3} \left[1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right) \right]}$$

...algebra...

$$\frac{\Gamma_2}{\Gamma_2 - 1} = \frac{\frac{5}{2} + D(y) \left[\frac{15}{4} + 5 \frac{\chi}{kT} + \left(\frac{\chi}{kT} \right)^2 \right]}{1 + D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right)}$$

$$D(y) = \frac{y(1-y)}{(1+y)(2-y)}$$

Similarly:

$$\Gamma_3 - 1 = \frac{2 + 2D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right)}{3 + 2D(y) \left(\frac{3}{2} + \frac{\chi}{kT} \right)^2}$$

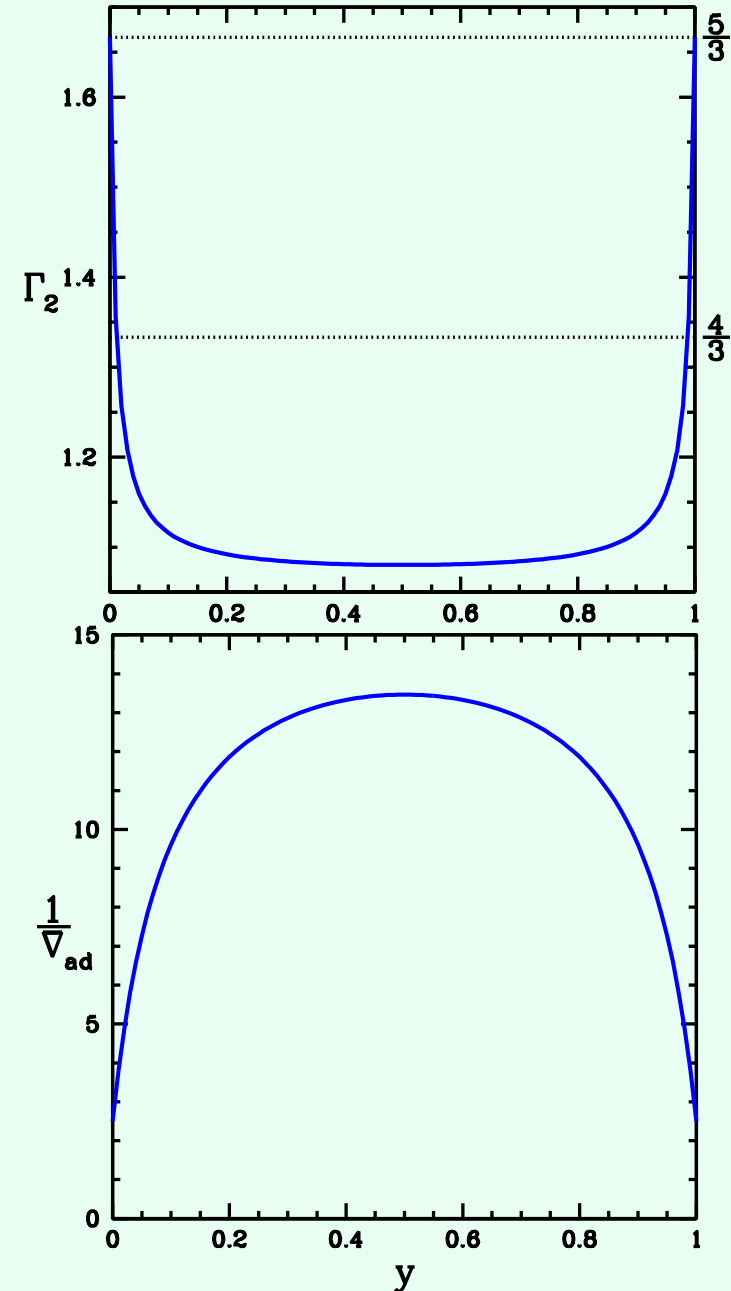
Adiabatic Exponents: Ionization Zones

$$\frac{\Gamma_2}{\Gamma_2 - 1} = \frac{1}{\nabla_{\text{ad}}} = \left(\frac{\partial \ln P}{\partial \ln T} \right)_{\text{ad}}$$

Gas heats up

- rapid increase of ionization
- increase in number of free particles
- pressure goes up faster than $\sim T$

Pressure increases faster than normal
In response to temperature increase



Adiabatic Exponents: Ionization Zones

$$\Gamma_3 - 1 = \left(\frac{\partial \ln T}{\partial \ln \rho} \right)_{\text{ad}}$$

Gas is compressed

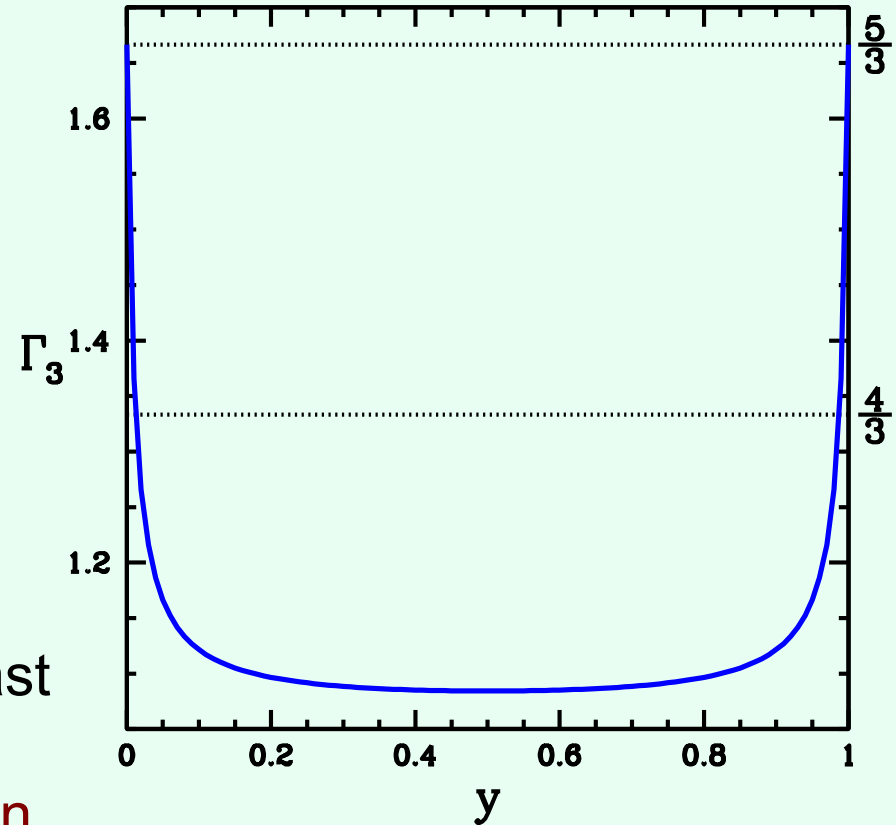
→ increase in temperature

→ rapid increase in ionization

→ energy goes into ionizing gas
instead of thermal energy

→ temperature does not rise as fast

Temperature increases slower than
normal in response to compression



Adiabatic Exponents: Ionization Zones

Adiabatic exponents can be less than $4/3 \rightarrow$ unstable!

Ionization zones are responsible for stellar pulsations

- second ionization of *He* provides the most significant driving
- ionization zone must be in the radiative region

Cool, luminous stars ($T_{\text{eff}} < 6300\text{K}$) will pulsate in an “instability strip” a few hundred K wide on HR diagram

